

Hidden Markov Models: Classical, Quantum, and Beyond

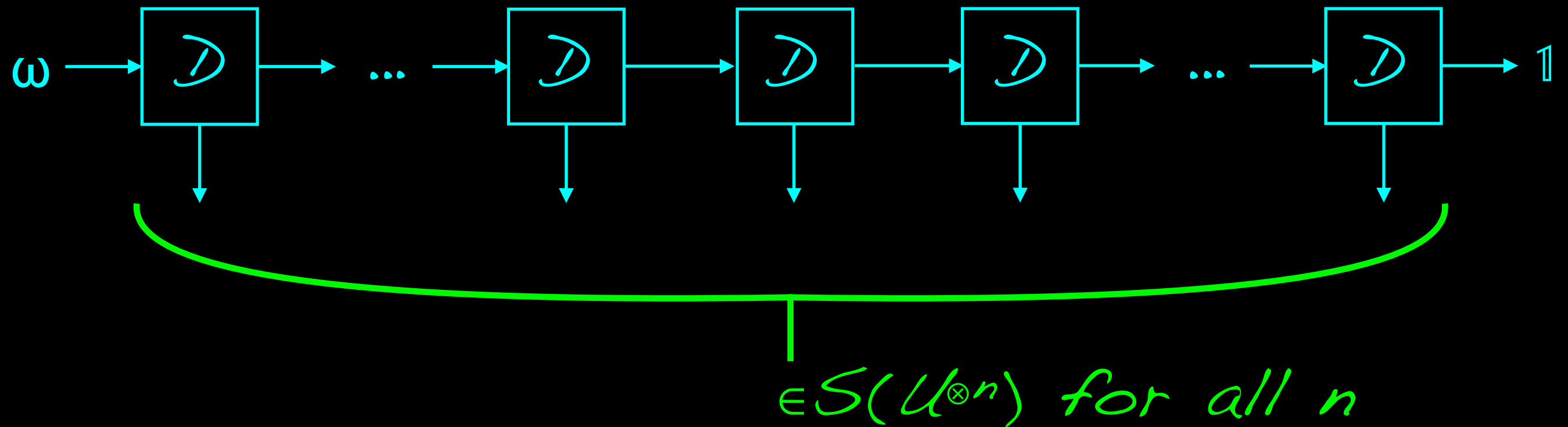
$$P(uv...w)$$



Andreas Winter
(ICREA & UAB Barcelona)

[A. Monràs/AW, JMP 2016 - 1412.3634;
M. Fanizza/J. Lumbreras/AW, CMP 2024 - 2209.11225]

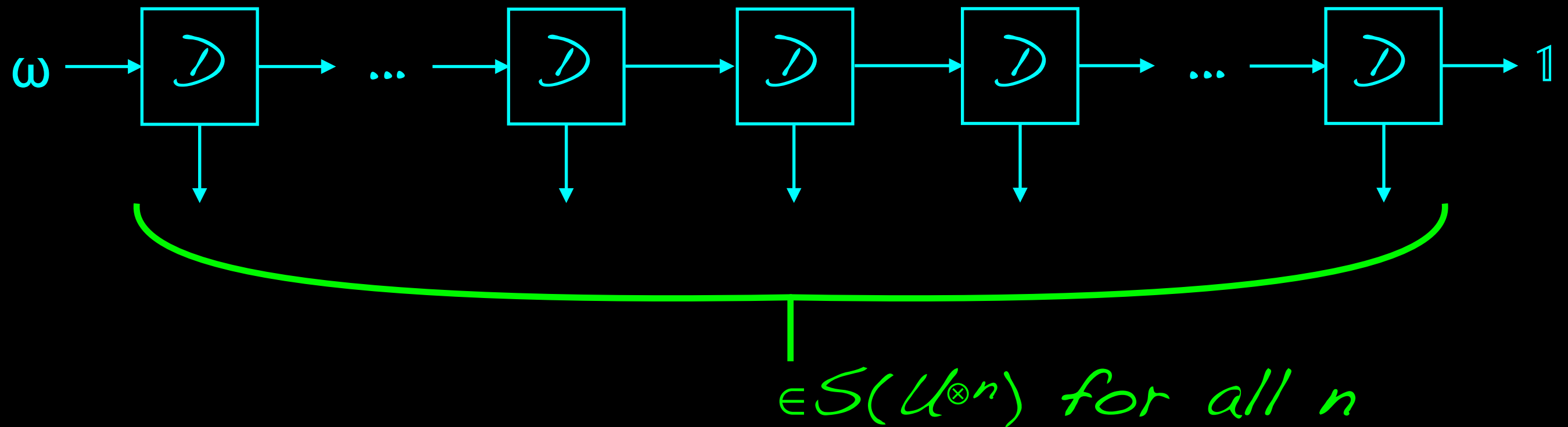
Finely correlated state given by cptp map $\mathcal{D}:S(\mathcal{H})\rightarrow S(\mathcal{H}\otimes\mathcal{U})$ and state $\omega = \text{Tr}_{\mathcal{U}}\circ\mathcal{D}(\omega)$:



When \mathcal{D} is (conjugation by) an isometry, we get matrix product state (MPS).

[Fannes/Nachtergaele/Werner, CMP 144:443-490 (1992)]

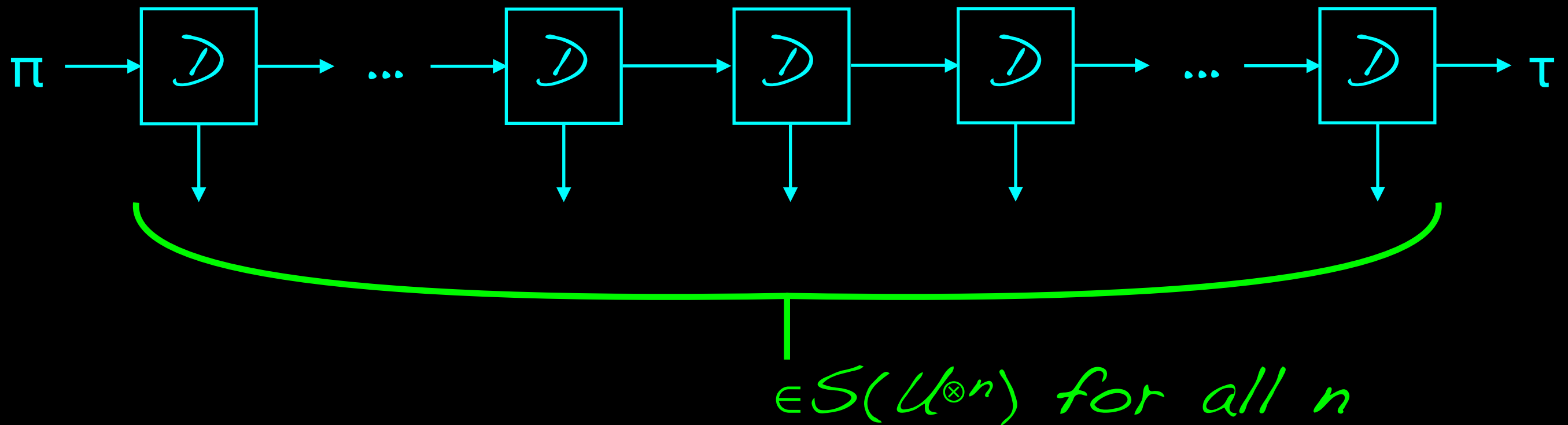
C^* -finitely correlated state given by cptp map $\mathcal{D}: S(\mathcal{H}) \rightarrow S(\mathcal{H} \otimes \mathcal{U})$ and state $\omega = \text{Tr}_{\mathcal{U}} \circ \mathcal{D}(\omega)$:



When \mathcal{D} is (conjugation by) an isometry, we get matrix product state (MPS).

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General finitely correlated state given by map
 $\mathcal{D}:V \rightarrow V \otimes B(\mathcal{U})$ and $V \ni \pi = \text{Tr}_{\mathcal{U}} \circ \mathcal{D}(\pi)$:



V is a vector space, τ linear s.t. $\tau(\pi)=1$.

What does this added generality actually buy beyond C^* -FCS?

[Fannes/Nachtergaele/Werner, CMP 144:443-490 (1992)]

Rest of the talk:

Classical finitely correlated states,
i.e. probability distribution on \mathbb{U}^∞

Concretely, we observe an infinite time
series $\dots u_{-k} \dots u_{-1} u_0 u_1 u_2 \dots u_k \dots$

[$u_t \in \mathbb{U}$ letters from a finite alphabet].

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Assume *stationarity*, i.e. for all t and ℓ ,

$$\Pr\{u_1 = u_1, \dots, u_\ell = u_\ell\} = \Pr\{u_t = u_1, \dots, u_{t+\ell-1} = u_\ell\}.$$

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$$\Pr\{u_1=u_1, \dots, u_\ell=u_\ell\} = \Pr\{u_t=u_1, \dots, u_{t+\ell-1}=u_\ell\}.$$

These marginals $P(\underline{u})$, for all finite words

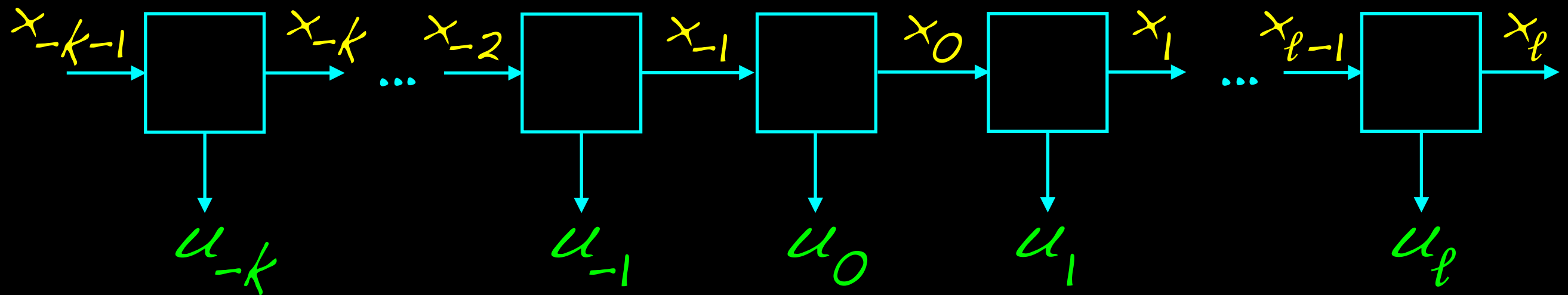
$$\underline{u} = u_1 u_2 \dots u_\ell \in \mathbb{U}^* = \bigcup_{k \geq 0} \mathbb{U}^k,$$

determine the probability law.

Rest of the talk:

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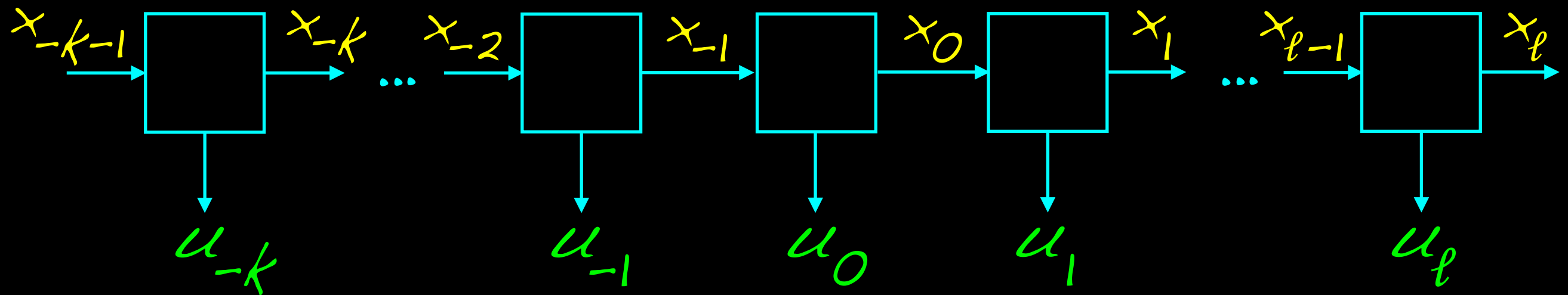
"Explanation" of $P(\underline{u})$ via a *finite memory system* as hidden cause:



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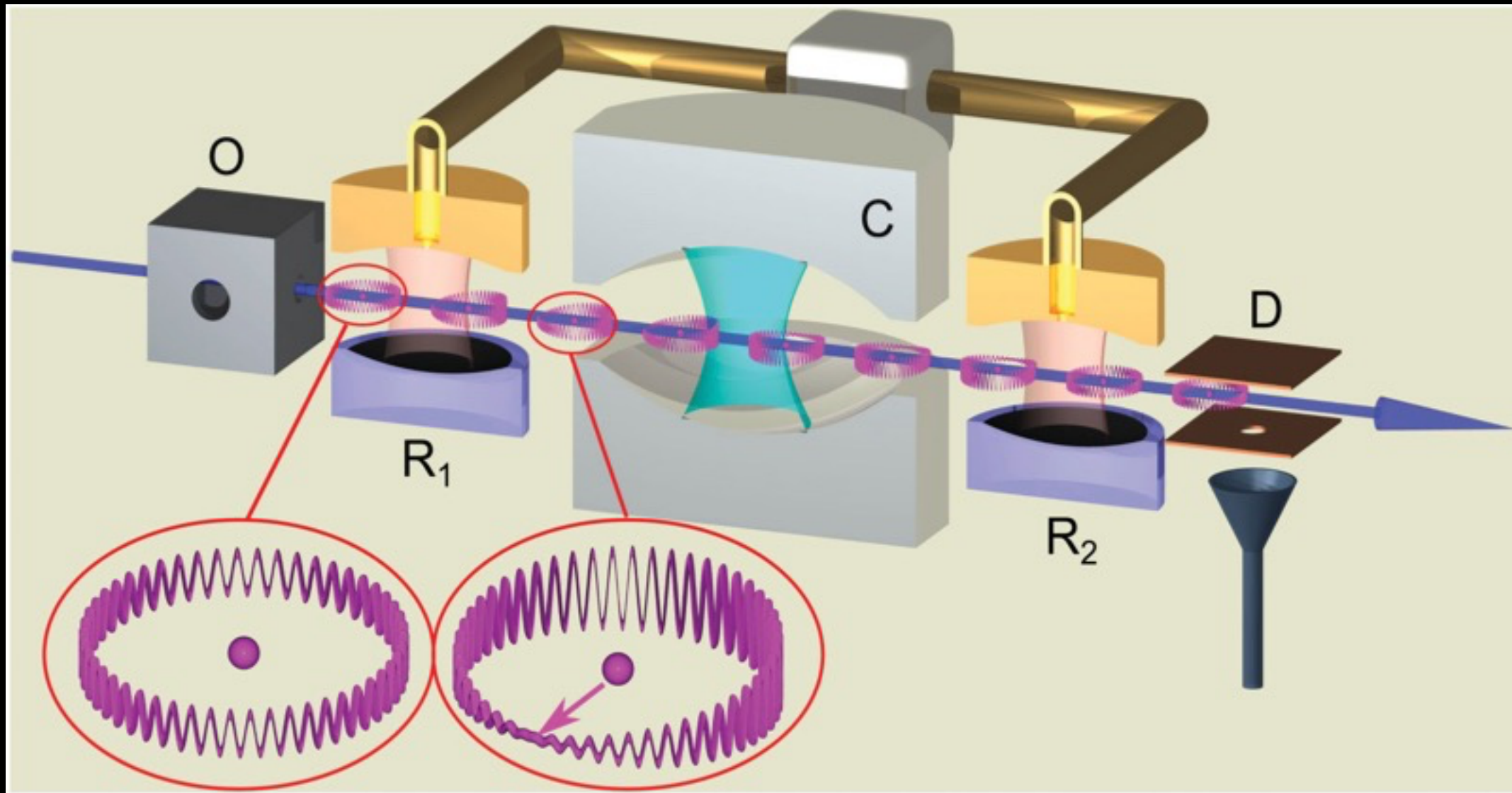
"Explanation" of $P(\underline{u})$ via a **finite memory system** as hidden cause:



Of course, need to specify the nature of the **causation**, and of the **memory**...

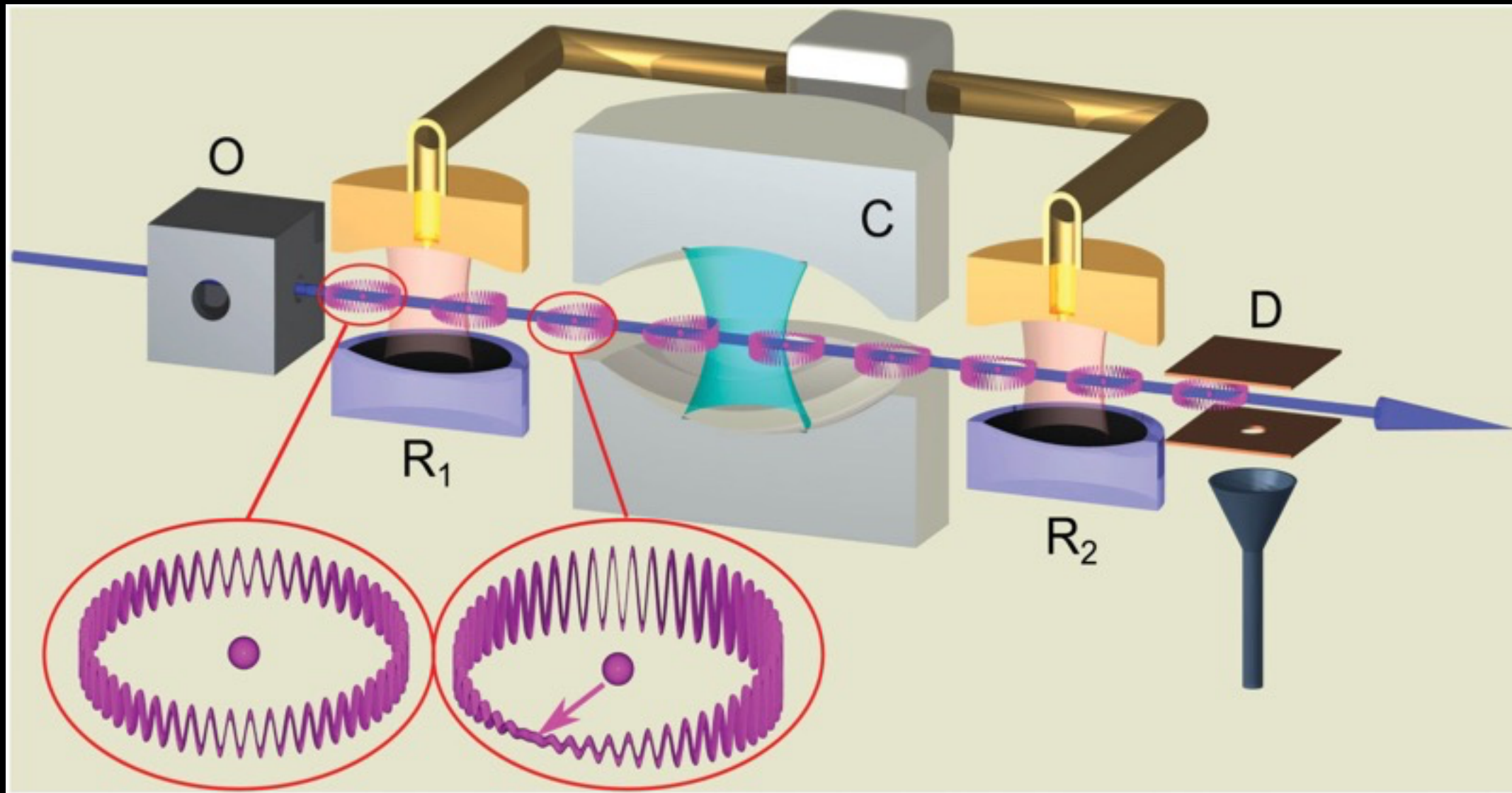
Example: Cavity-atom interaction

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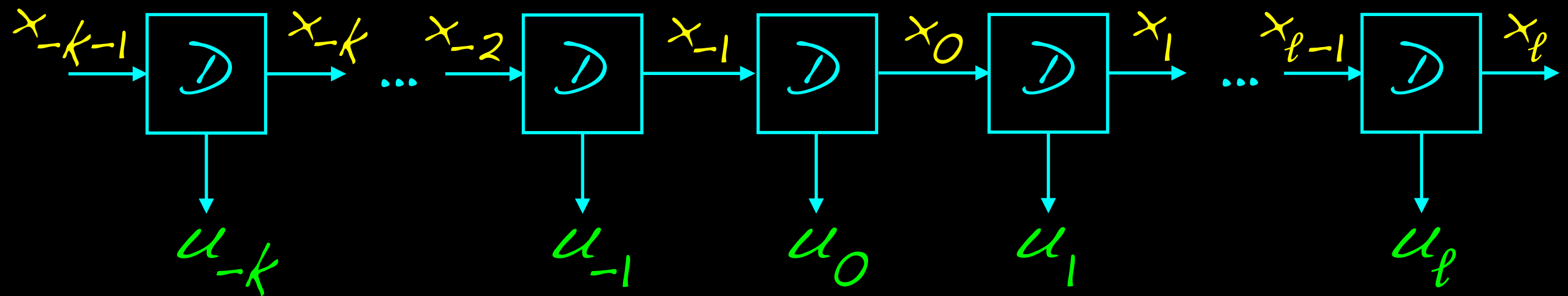
Question: can one infer the quantum nature of the internal mechanism by observing $P(\underline{u})$?

Outline

- ✓. Observations as consequence of a finitary hidden cause (memory)
 - 1. Classical, quantum and GPT memory
 - 2. Reconstructing a quasirealisation:
low-rank Hankel matrix (completion)
 - 3. Separations: classical $\stackrel{\checkmark}{\subsetneq}$ quantum $\stackrel{\checkmark}{\subsetneq}$ GPT

1-a. Classical memory (HMM)

The $x_t \in \mathbb{X}$ are from a finite set of internal states, $\mathcal{D}: \mathbb{X} \rightarrow \mathbb{X} \times \mathbb{U}$ is a stochastic map:



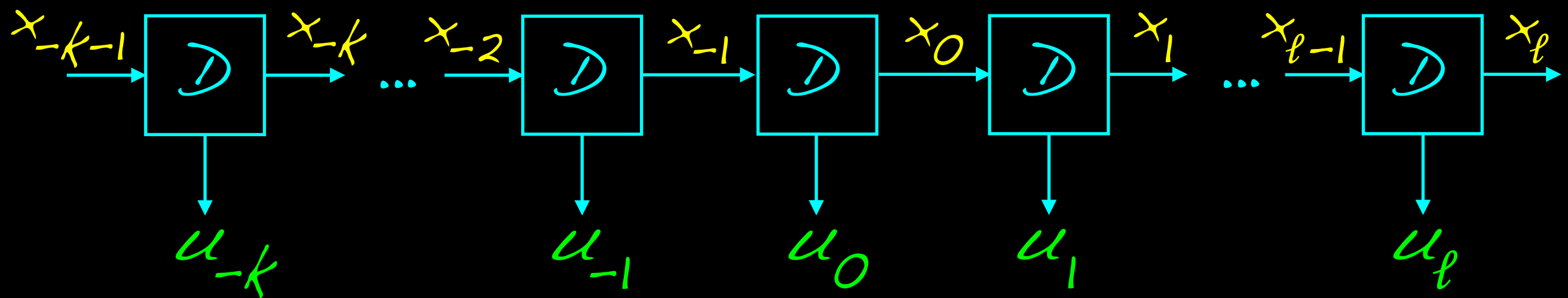
$\mathcal{D}_u: \mathbb{X} \rightarrow \mathbb{X}$ are sub-stochastic maps, s.t.

$\bar{\mathcal{D}} = \sum_u \mathcal{D}_u$ is stochastic with stationary distribution π : $\bar{\mathcal{D}}\vec{1} = \vec{1}$, $\pi\bar{\mathcal{D}} = \pi$.

$$P(u_1 u_2 \dots u_\ell) = \pi \mathcal{D}_{u_1} \mathcal{D}_{u_2} \dots \mathcal{D}_{u_\ell} \vec{1} \quad (\text{p.r.})$$

1-6. Quantum memory (\mathcal{H} /QMM)

The $x_t \in \mathbb{X} = S(\mathcal{H})$ are quantum states on \mathcal{H} ,
and \mathcal{D} is a completely positive instrument:



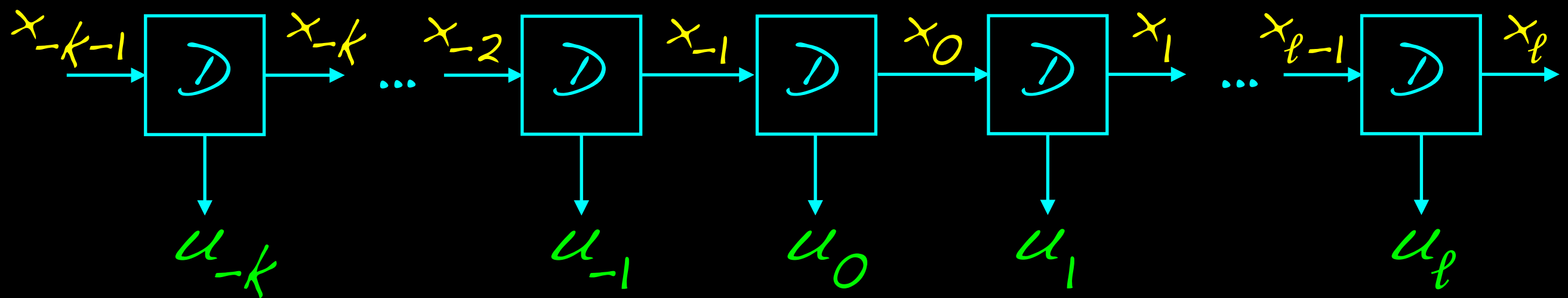
$\mathcal{D}_u : \mathbb{X} \rightarrow \mathbb{X}$ are completely positive maps, s.t.

$\bar{\mathcal{D}} = \sum_u \mathcal{D}_u$ is unital (cpup) with stationary state $\omega : \bar{\mathcal{D}}\mathbb{1} = \mathbb{1}, \omega \circ \bar{\mathcal{D}} = \omega$.

$$P(u_1 u_2 \dots u_\ell) = \omega \circ \mathcal{D}_{u_1} \circ \mathcal{D}_{u_2} \dots \circ \mathcal{D}_{u_\ell} \mathbb{1} \quad (\text{c.p.r.})$$

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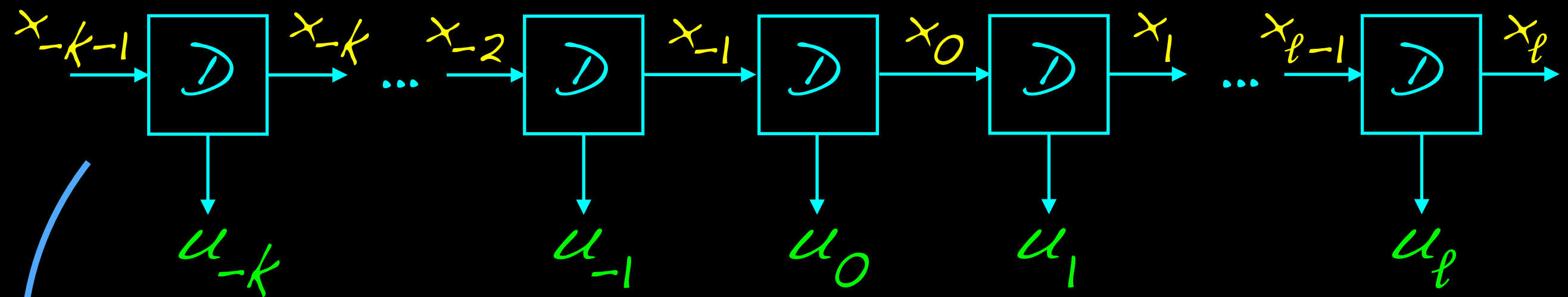
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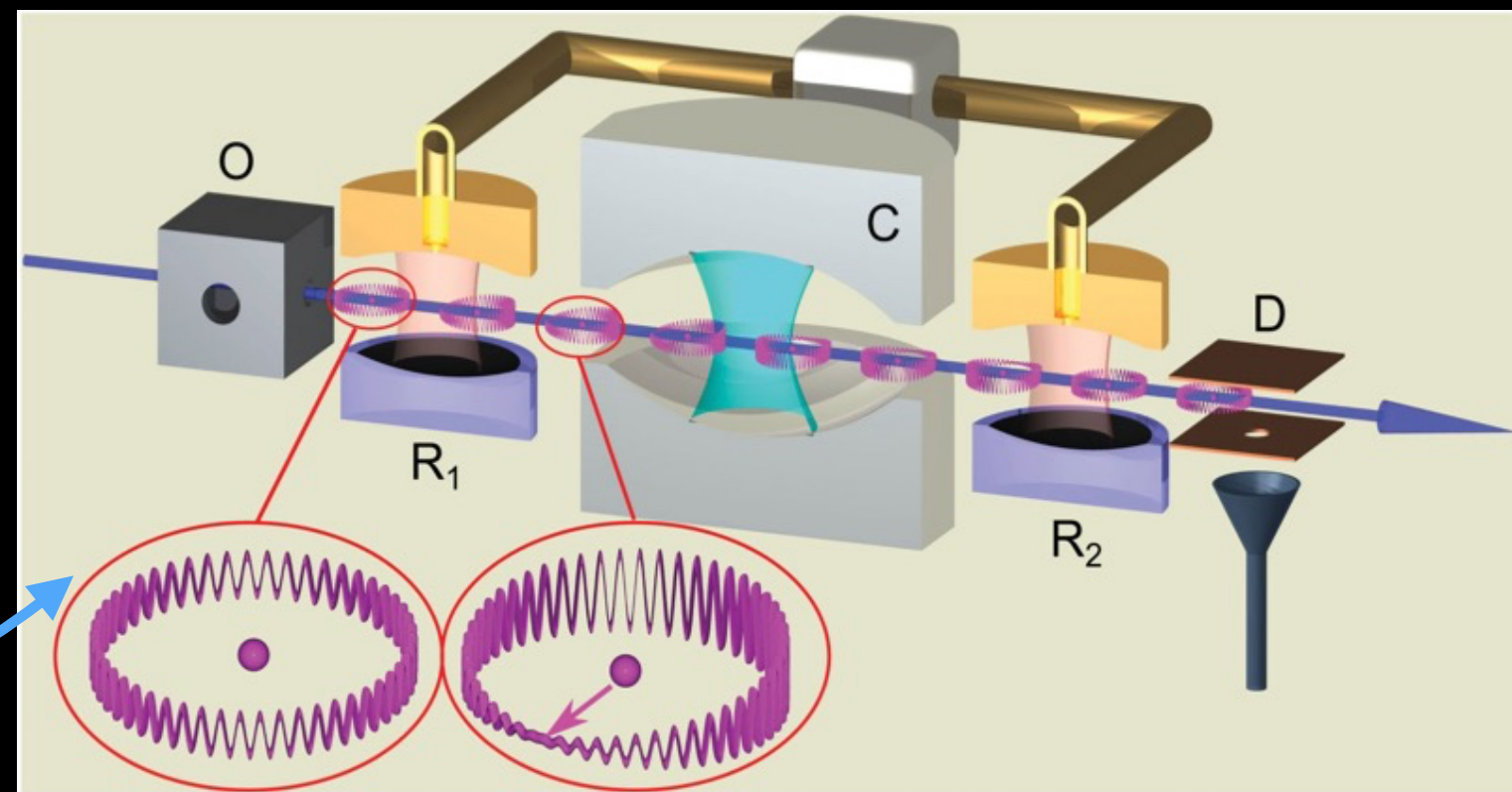
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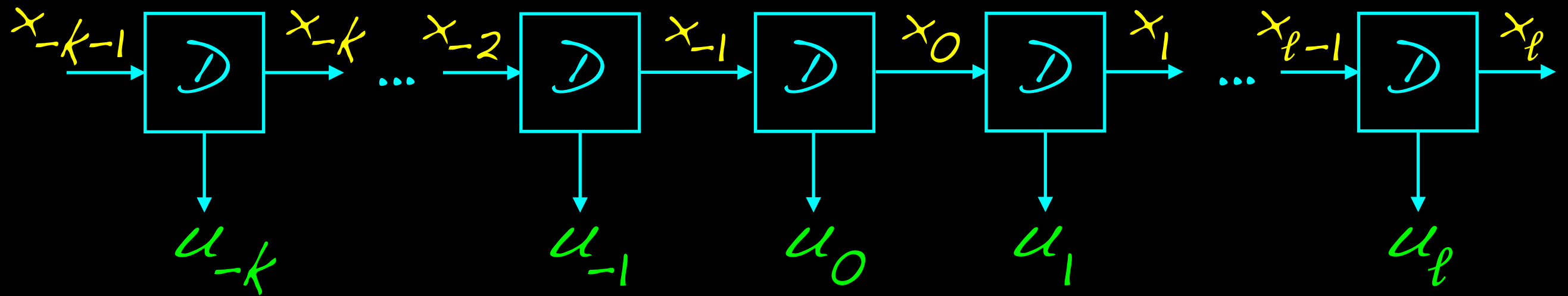


In real life (=in
the laboratory):



1-c. General linear structure

The $x_t \in V$ are elements of a (real) vector space, and \mathcal{D} is a collection of linear maps:



$\mathcal{D}_u : V \rightarrow V$ are linear maps, $\tau \in V$, $\pi \in V^*$, s.t.

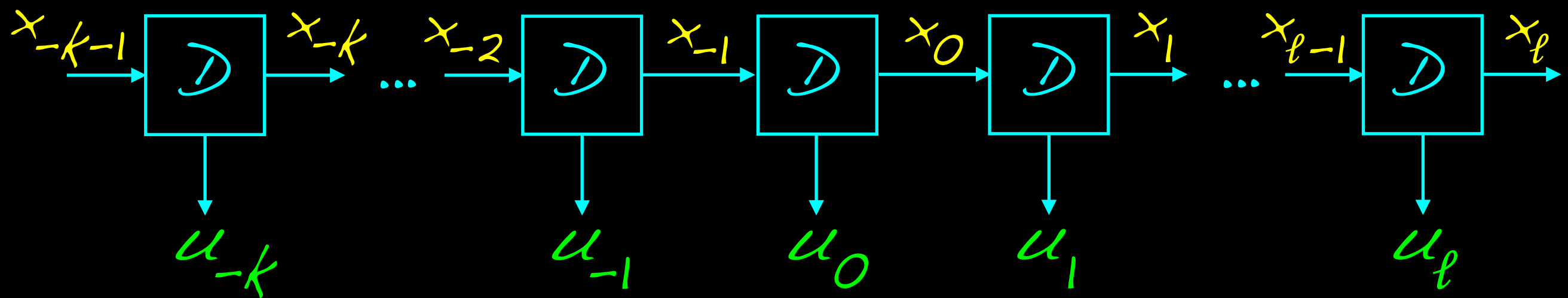
$\bar{\mathcal{D}} = \sum_u \mathcal{D}_u$ preserves both τ and ω :

$\bar{\mathcal{D}}\tau = \tau$, $\pi \circ \bar{\mathcal{D}} = \pi$, as well as $\pi(\tau) = 1$.

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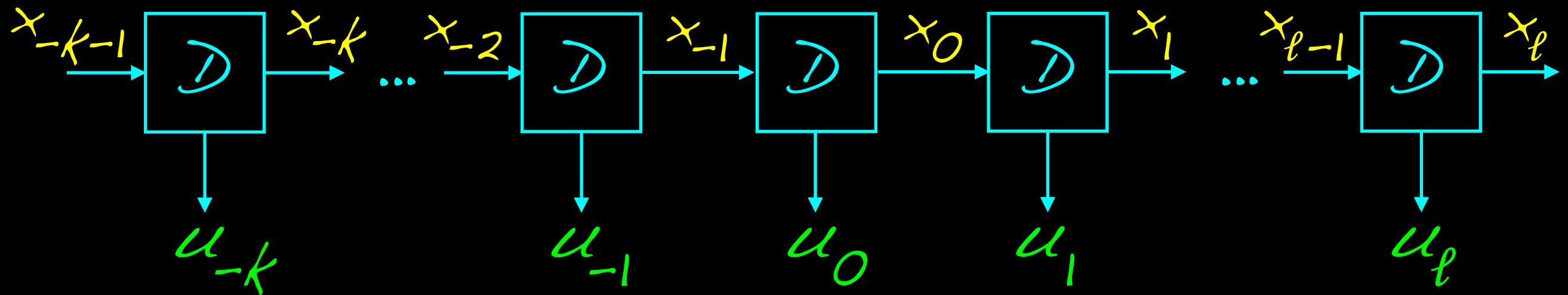
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Quasirealisation:
 general finitely
 correlated state

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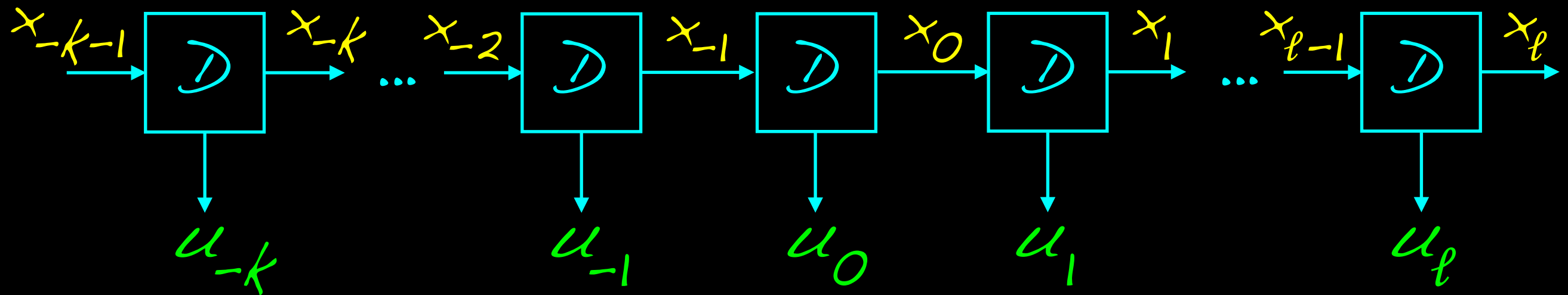
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Unlike classical and quantum case, no a priori guarantee that $P(\underline{u}) \geq 0$. In fact, checking positivity is undecidable ⚡

[Sontag, J. Comp. Syst. Sci. 11(3):375-381, 1975]

Simple examples: - I.i.d. distributed $u_t \in \mathbb{U}$,
i.e. $P = P_1^{\otimes \mathbb{Z}}$ infinite product of single-letter
distributions P_1 . Requires no, or rather only
trivial, memory: $\dim V = 1$.

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-De Finetti distribution $P = \sum_{x \in \mathbb{X}} \pi_x P_x^{\otimes \mathbb{Z}}$ with
distinct p.d.'s P_x and $\pi_x > 0$.

Realised as HMM by memorising $x \in \mathbb{X}$ forever,
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For $|\mathbb{X}| = \infty$: stationary process not realised as
HMM, HQMM, or even quasirealisation.

Example. $V = B(\mathbb{C}^2)_{sa} = \text{span}\{1, X, Y, Z\}$ qubit
with $\tau = 1$, $\pi = \frac{1}{2}\text{Tr}$, and the following maps:

$$D_0(A) = \frac{1}{4} |0\rangle\langle 0| A |0\rangle\langle 0|,$$

$$D_1(A) = \frac{1}{4} |1\rangle\langle 1| A |1\rangle\langle 1|,$$

$$D_X(A) = \frac{1}{4} \exp(i\alpha X) A \exp(-i\alpha X),$$

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When α/π and β/π are irrational, dynamics
 explores whole Bloch sphere densely. Four-
 dim. q.u.r., but requires 2 qubits for c.p.r.!

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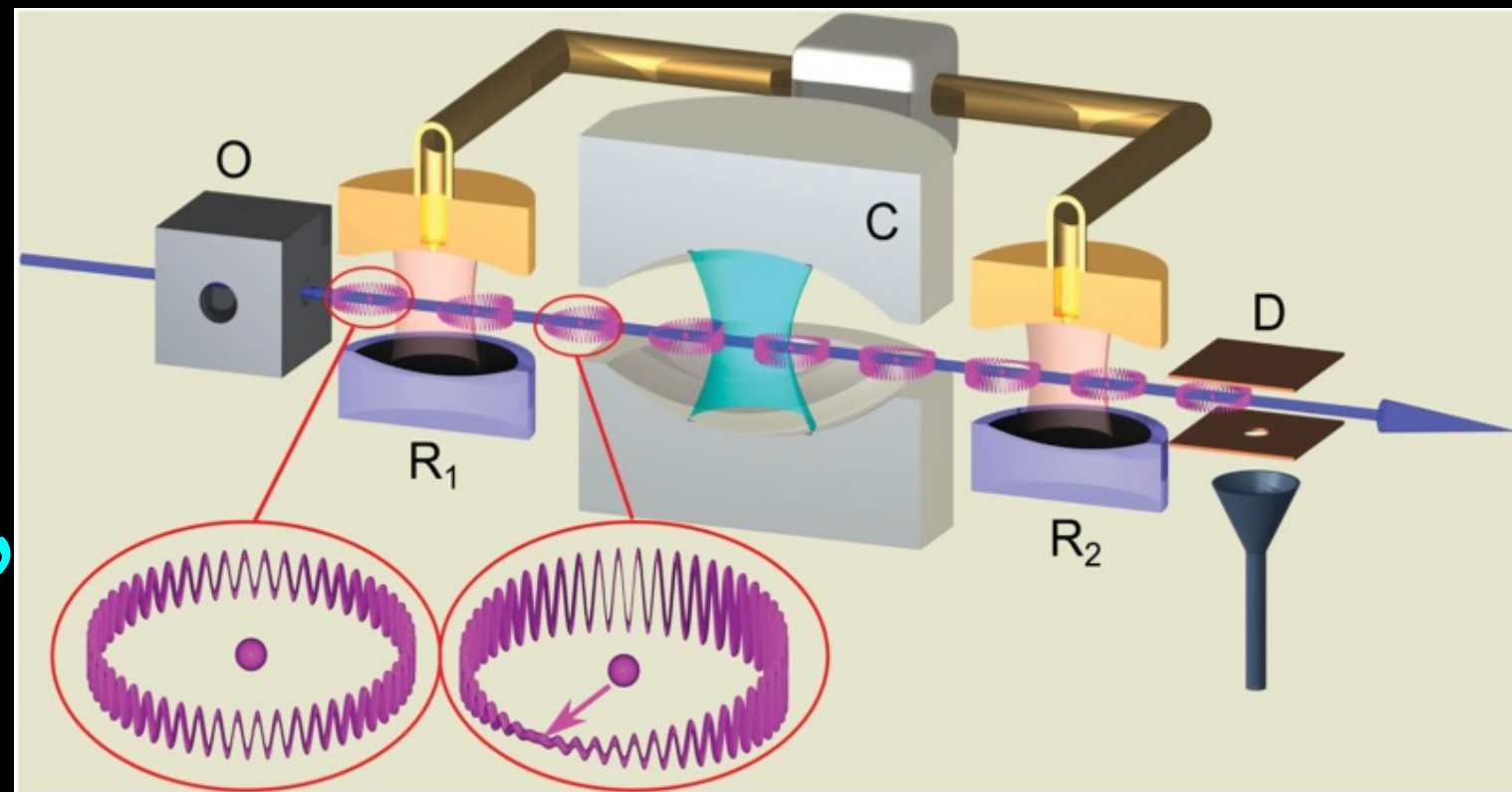
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explores whole Bloch sphere densely. Four-
dim. q.u.r.: $\mathcal{H}QMM$ with qubit memory.

Recover the internal mechanism from $P(\underline{u})$?

Quantum application: characterisation of quantum devices - state preparation, gates and measurements - from first principles.

[R. Blume-Kohout et al., 1310.4492] treat system as a black box whose reaction to different interventions we can observe...

Evidently possible only up to linear equivalence, e.g. isometries.



What guarantees positivity of probability?

* $\underline{u} = u_1 u_2 \dots u_\ell \mapsto \mathcal{D}_{\underline{u}} = \mathcal{D}_{u_1} \circ \mathcal{D}_{u_2} \dots \circ \mathcal{D}_{u_\ell}$ is semigroup representation.

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* Classical & quantum case: positivity $P(\underline{u}) \geq 0$ enforced by the vector space order.

Generally: Assume we have convex cones $C \subset V$ and $\tilde{C} \subset C' \subset V^*$, s.t. $\tau \in C$, $\pi \in \tilde{C}$, and the cones are preserved by the transformations, i.e. $D_u C \subset C$, $\tilde{C} D_u \subset \tilde{C} \forall u$. Then $P(\underline{u}) \geq 0$.

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Dual cone $C' = \{f \in V^* : f(x) \geq 0 \forall x \in C\}$

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Conversely: If $P \geq 0$, then such cones exist,

e.g. $C = C_{\min} = \overline{\text{cone}\{D_{\underline{u}} \tau : \underline{u} \in U^*\}},$

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...not unique, could for instance also take dual cone $\tilde{C} = C'$; call any such C "suitable".

Interpretation: finite-dimensional
quasirealisation "explains" time series
 P by the hidden mechanism of a
general probabilistic theory (GPT):

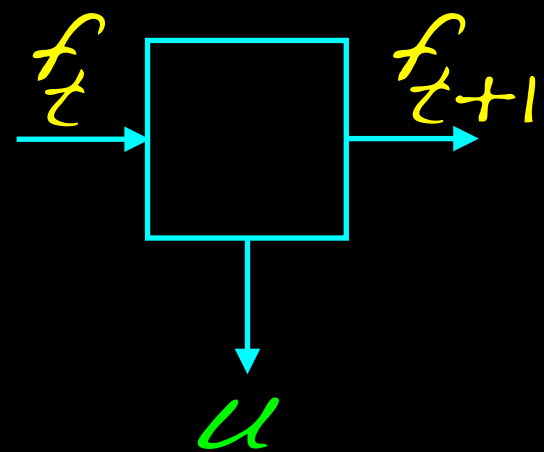
- C and C' are pointed and generating cones;
- $\tau \in \text{int}(C)$ and $S := \{f \in C' : f(\tau) = 1\}$ state space;
- $\mathcal{E} := C_n(\tau - C)$ "effects" for measurements.

[G. Ludwig & school, 1960s-70s, ...]

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$\equiv \left\{ \begin{array}{l} f_t \circ D_u = \Pr\{u | f_t\} f_{t+1}, \text{ relates} \\ \text{current \& future states,} \\ \text{and the output } u. \end{array} \right.$

2. Reconstruction of V

* Consider the Hankel-type matrix $\mathcal{H} = (\mathcal{H}_{\underline{u}, \underline{v}})$,
with $\mathcal{H}_{\underline{u}, \underline{v}} = \mathcal{P}(\underline{u}\underline{v}) = \mathcal{P}(u_1 u_2 \dots u_\ell v_1 v_2 \dots v_k)$
 $= \mathcal{H}_{\varepsilon, \underline{u}\underline{v}} = \mathcal{H}_{\underline{u}\underline{v}, \varepsilon}$.

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* If the process \mathcal{P} has a quasirealisation of
 $\dim V = d$, then

$$\mathcal{H}_{\underline{u}, \underline{v}} = (\pi^\circ \mathcal{D}_{\underline{u}})(\mathcal{D}_{\underline{v}}^\top),$$

and so $\text{rank } \mathcal{H} \leq d$.

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Thus: finite rank of \mathcal{H} necessary requirement for the existence of a quasirealisation, and hence of classical or quantum hidden Markov models.

Is it sufficient?

* Consider the Hankel-type matrix $\mathcal{H} = (\mathcal{H}_{\underline{u}, \underline{v}})$,
with $\mathcal{H}_{\underline{u}, \underline{v}} = P(\underline{u}\underline{v}) = P(u_1 u_2 \dots u_\ell v_1 v_2 \dots v_k)$.
Necessarily of finite rank.

* Conversely, if $\text{rank } \mathcal{H} = r < \infty$: There exists
a g.u.r. ("regular rep.") with $\dim V = r$,
which is the minimum. Any other minimal-
dim. g.u.r. is similar to the regular one,
i.e. linearly equivalent.

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[Construction: $V =$ column space of \mathcal{H} , and
 D_u maps $h_{\underline{v}} = \mathcal{H}_{\cdot, \underline{v}}$ to $h_{u\underline{v}} = \mathcal{H}_{\cdot, u\underline{v}} = \mathcal{H}_{\cdot, \underline{u}, \underline{v}}$ -
 linear because it selects the rows of $h_{\underline{v}}$
 with index ending in u ; $\tau = h_\varepsilon$, $\pi = (1, 0, 0, \dots)$.
 Check that it works...]

Fine, so assume finite rank r of $\mathcal{H} = (\mathcal{H}_{\underline{u}, \underline{v}})$,
i.e. an r -dimensional quasirealisation exists.

Is the process then generated by a finite
memory $\mathcal{H}MM$ (classical p.r.)?

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NO! [Fox/Rubin (1968) and Dharmadhikari/
Nadkarni (1970)] provided first examples of
processes with finite Hankel rank (actually
 $r=3$) but requiring infinite classical memory.
In fact they're defined as $\mathcal{H}MM$ w/ infinite
memory. Exploits spectral information from
Perron-Frobenius theory.

Alright: assume finite rank r of $\mathcal{H}=(\mathcal{H}_{\underline{u},\underline{v}})$,
i.e. an r -dimensional quasirealisation exists.

Is the process then generated by a finite
memory \mathcal{H} QMM (quantum c.p.r.)?

(Asked by Fannes/Nachtergaele/Werner

[CMP 144:443-490 (1992)] for general finitely
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Find: Fox/Rubin's and Dharmadhikari/
Nadkarni's processes have \mathcal{H} QMMs (with
gutruts).

[M. Fanizza/J. Lumbrellas/AW, arXiv:2209.11225]

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Minimal-dimensional quasirealisation of a process is unique, and isomorphic to the regular representation from \mathcal{H} , $\dim V = r$.

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Fact: Given any quasirealisation V , then the regular one is obtained by going to quotient $V_0 =: \text{span}(C_{\min}) / \ker(C'_{\max})$.

For the cone C (classical, quantum or GPT), this means intersecting it with $\text{span}(C_{\min})$, and factoring out $\ker(C'_{\max})$.

Recall cones:

Given convex cones $C \subset V$ and $\tilde{C} \subset C' \subset V^*$,
s.t. $\tau \in C$, $\pi \in C'$, and the cones are preserved by
the transformations, i.e. $D_u C \subset C$, $\tilde{C} D_u \subset \tilde{C}$
for all u . Then $P(\underline{u}) \geq 0$.

Conversely: If $P \geq 0$, then such cones exist,

$$\text{e.g. } C = C_{\min} = \overline{\text{cone}\{D_u \tau : \underline{u} \in U^*\}},$$

$$\tilde{C} = C'_{\max} = \overline{\text{cone}\{\pi D_u : \underline{u} \in U^*\}}.$$

But not unique: many cones between C_{\min}
and C_{\max} are suitable: $C_{\min} \subset C \subset C_{\max}$.

(Also, of course, C has to be stable under
the maps D_u !)

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A $\mathcal{H}MM$ (p.r.) has the cone of non-negative vectors; this gives rise to *polyhedral cones* C & C' in the regular representation.

A $\mathcal{H}QMM$ (c.p.r.) has cone of semidefinite matrices; this gives rise to *semidefinite representable (SDR) cones* C & C' in the regular representation:

$$C = \{x = (x_1, \dots, x_d) : \exists x_{d+1}, \dots, x_e \sum_{i=1}^e x_i R_i \geq 0\},$$

for certain $D \times D$ -matrices R_i .

Example. $V = B(\mathbb{C}^2)_{sa} = \text{span}\{1, X, Y, Z\}$ qubit
with $\tau=1$, $\pi=\frac{1}{2}\text{Tr}$, and the following maps:

$$\mathcal{D}_0(A) = \frac{1}{4} |0\rangle\langle 0| A |0\rangle\langle 0|,$$

$$\mathcal{D}_1(A) = \frac{1}{4} |1\rangle\langle 1| A |1\rangle\langle 1|,$$

$$\mathcal{D}_X(A) = \frac{1}{4} \exp(i\alpha X) A \exp(-i\alpha X),$$

$$\mathcal{D}_Z(A) = \frac{1}{4} \exp(i\beta Z) A \exp(-i\beta Z),$$

$$\mathcal{D}_T(A) = \frac{1}{4} A^T.$$

When α/π and β/π are irrational, dynamics explores whole Bloch sphere densely. Four-dim. q.u.r., but requires 2 qubits for c.p.r.!

In the previous example,

$C_{\min} = C_{\max}$, hence C is

unique, and it's not polyhedral: cone over Bloch sphere. Thus, this process has no (finite) classical realisation.

Polyhedral cone between C_{\min} and C_{\max} necessary for cl. realisation. Sufficient?

[Cf. however Dharmadhikari/Nadkarni]

SDR cone between C_{\min} and C_{\max} necessary
for existence of a quantum realisation.

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Thm. [M. Fanizza/J. Lumbraeras/AW, 2209.11225]:
 \exists process P with Hankel rank $H = 3$ and
 $C_{\min} = C_{\max}$ transcendental, whereas SDR
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SDR cone between C_{\min} and C_{\max} necessary for existence of a quantum realisation.
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Thm. [M. Fanizza/J. Lumbraeras/AW, 2209.11225]:
 \exists process P with Hankel rank $\mathcal{H} = 3$ and $C_{\min} = C_{\max}$ transcendental, whereas SDR cones are semi-algebraic. Thus, P has no \mathcal{H} QMM.

Answers open question of Fannes/Nachtergaele/Werner [CMP 144:443-490 (1992)] :-)

Example: P has three symbols, 0, 1, 2.

We give directly its quasirealisation:

$V = \mathbb{R}^3$; let $a > 1 > b > 0$ such that $\ln a$ and $\ln b$ are linearly independent over the rationals.

$$\mathcal{D}_1 = \begin{bmatrix} a & 0 & 0 \\ 0 & 1 & 0 \\ 0 & \ln a & 1 \end{bmatrix}, \mathcal{D}_2 = \begin{bmatrix} b & 0 & 0 \\ 0 & 1 & 0 \\ 0 & \ln b & 1 \end{bmatrix},$$

$$\mathcal{D}_0 = m_0 \mu_0^T, \text{ with } m_0 = \begin{bmatrix} m_{01} \\ m_{02} \\ m_{03} \end{bmatrix},$$
$$\mu_0^T = [\mu_{01} \ \mu_{02} \ \mu_{03}]$$

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$$\mathcal{D}_1^s \mathcal{D}_2^t = \begin{bmatrix} a^s b^t & 0 & 0 \\ 0 & 1 & 0 \\ 0 & s(\ln a) + t(\ln b) & 1 \end{bmatrix} \text{ for } s, t \in \mathbb{N}$$

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$$\mathcal{D}_1^s \mathcal{D}_2^t = \begin{bmatrix} e^x & 0 & 0 \\ 0 & 1 & 0 \\ 0 & x & 1 \end{bmatrix}, \text{ with } x \in \mathbb{R} \text{ dense!}$$

Example (cont'd): D_0 is a "reset" operation (making P a "birth process"), so can write

$$C'_{\max} = \text{cone}\{\mu_0 e^x \quad \mu_0 + \mu_0 x \quad \mu_0\} : x \in \mathbb{R}\}$$

$$C_{\min} = \text{cone}\{m_0 e^x \quad m_0 \quad m_0 + m_0 x\}^T : x \in \mathbb{R}\}$$

Fact: C_{\max} is of the same form as C_{\min} , only with different parameters. One can choose D_0 such that $C_{\min} = C_{\max} = C$.

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Fact: C_{\max} is of the same form as C_{\min} , only with different parameters. One can choose D_0 such that $C_{\min} = C_{\max} = C$.

In that case, a suitable positive linear combination of D_0, D_1, D_2 has right fixed point τ in $\text{int}(C)$, and left fixed point π in $\text{int}(C')$. This is the sought-after g.u.r. (...)

Example gives rise to the exponential cone

$$K_{\text{exp}} = \{(x, y, z) : x/y \geq e^z/y, x, y, z \geq 0\},$$

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More examples from power cone ($0 < t < 1$)

$$K_t = \{(x, y, z) : x^t y^{1-t} \geq |z|, x, y \geq 0, z \in \mathbb{R}\},$$

which is transcendental iff t is irrational.

As before we can design a reset map and two invertible maps, which latter act densely transitive on the cone's extremal rays. And we can engineer $C_{\min} = C_{\max}$, too.

4. Further thoughts

- Exhibited a process (FCS) without a quantum realisation (i.e. it is not C^* -FCS). However, that is an asymptotic statement, every finite block of the chain is c.p. representable. To approximate P on n sites to error ε , how large a virtual dimension t do we need? For fixed t and $n \rightarrow \infty$, does ε become arbitrarily small, or is it bounded away from 0, or does even converges to 1?

4. Further thoughts

Extend to genuinely quantum case, i.e. a chain of non-commutative spin C^* -algebras:

- Have a generalisation of regular (minimum dim.) representation for *finitely corr. states*
- Rather than a vector space order on V and positive elements and maps, necessary and sufficient structure is an *operator system*, i.e. consistent orders on $V \otimes M_n$, and maps are completely positive...

[Fannes/Nachtergaele/Werner, CMP 144:443-490 (1992)]

4. Further thoughts

Finely correlated state on a chain of non-commutative spin C^* -algebras:

- In fact, the finely correlated state itself gives us two extreme o.s., where the cones $(V \otimes M_n)_+$ are either all as small or all as large as they can be.
- Exponential and power cones have matrix generalisations; perhaps suitable for new variational classes of FCS? Need cp maps!

[Fanizza et al., work in progress]

4. Further thoughts

Finitely correlated state on a chain of non-commutative spin C^* -algebras:

- Examples of a FCS that are not C^* -FCS are highly mixed (they're p.d.'s). So what about pure ones?
- Note that C^* -FCS always have C^* -FCS purifications. Do our example FCS have FCS purifications?

[Fanizza et al., work in progress]

4. Questions, questions, questions

- Low-rank completion of the Hankel matrix with noisy data? Cf. [Fanizza/Galke/Lumbreras/Rouzé/AW, arXiv:2312.07516]
- How to find a quantum model just from the regular representation (assuming one exists)?
- Can these exponential and power cones be useful? Note that dual cone is of the same kind, so perhaps good for convex optimisation. Interesting class of GPTs?



=Additional material=

Ogni scarrafon' è
bell' a mamma suja



5. Removing redundancy: quotients

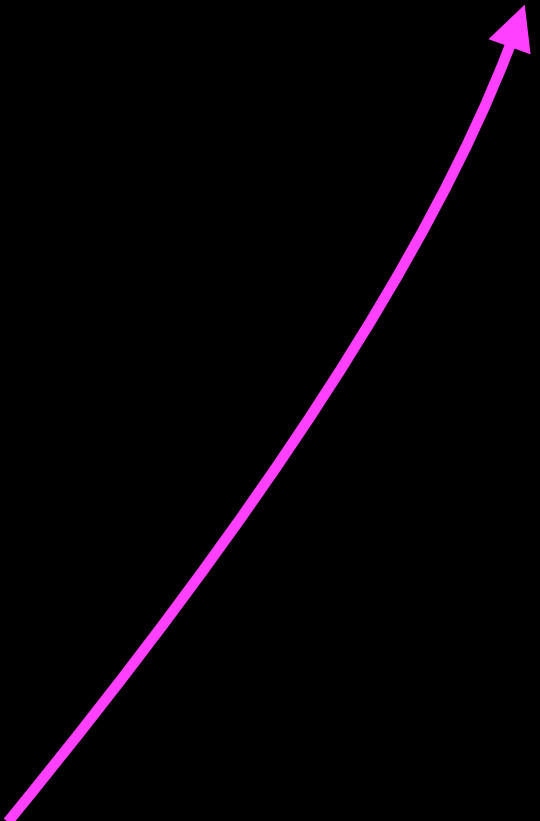
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assuming it has a suitable cone CcV .

Redundancy...

5. Removing redundancy: quotients

If your model is not minimal, still useful, assuming it has a suitable cone $C \subset V$.

Redundancy: $\mathcal{W} = \text{span}\{\mathcal{D}_{\underline{u}}^T : \underline{u} \in \mathbb{U}^*\} \subset V$,



Reachable space; might as well go to \mathcal{W} , with cone $C \cap \mathcal{W}$...

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$$K = \{\pi \mathcal{D}_{\underline{u}} : \underline{u} \in U^*\}^\perp \subset V.$$

Null space; $C \cap K = 0$, so we may factor out K ...

Reachable space; might as well go to \mathcal{W} , with cone $C \cap \mathcal{W}$...

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$$K = \{\pi \mathcal{D}_{\underline{u}} : \underline{u} \in \mathbb{U}^*\}^\perp \subset V.$$

$$V_0 := \mathcal{W}/K,$$

$$C_0 := (C \cap \mathcal{W})/K = \{\omega + K : \omega \in C \cap \mathcal{W}\},$$

$\tau_0 := \tau + K$, $\pi_0 := \pi/K$, $\mathcal{D}_{\underline{u}}^0 := \mathcal{D}_{\underline{u}}/K$; well-defined because of $\pi(K) = 0$, $\mathcal{D}_{\underline{u}}\mathcal{W} \subset \mathcal{W}$, $\mathcal{D}_{\underline{u}}K \subset K$.

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Always a minimal-dim. g.u.r., hence is isomorphic to regular, and cone C_0 is suitable.

Redundancy: $\mathcal{W} = \text{span}\{\mathcal{D}_{\underline{u}}^\top : \underline{u} \in \mathbb{U}^*\} \subset V$,

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Classical model, i.e. $V = \mathbb{R}^d$, $C = \mathbb{R}_{\geq 0}^d$, $\tau = (1, \dots, 1)^\top$,

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π a probability row vector.

C_0 is then a *polyhedral cone* and every such cone arises in the above way (Fourier-Motzkin elimination). Guaranteed:

$d \leq \# \text{extremal rays of } C$, sometimes best.

5'. Quotient of a quantum model

Quantum model, i.e. $V = B(\mathcal{H})_{sa}$, $C = B(\mathcal{H})_{\geq 0}$, $\tau = 1$,

$\pi = \omega$ quantum state, \mathcal{D}_u are cp maps.

Once constructed $K \cap \mathcal{W} \subset \mathcal{W} \subset V$: $C \cap \mathcal{W}$ is an operator system, $C_0 = (C \cap \mathcal{W})/K$ a quotient operator system; the \mathcal{D}_u^0 preserve C , in fact cp maps in the operator system sense.

[Farenick/Paulsen, Math. Scand. III:210-243, 2012]

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[Farenick/Paulsen, Math. Scand. III:210-243, 2012]

Membership in the cone is an SDP: semi-definite condition of a finite-size matrix with existential real variables.

SDR operator systems:

$$\mathbb{1} \in \mathcal{W} = \text{span}\{\mathcal{D}_{\underline{u}}\mathbb{1} : \underline{u} \in \mathbb{U}^*\} = \mathcal{B}(\mathcal{H})_{sa},$$

$$K = \{\omega \circ \mathcal{D}_{\underline{u}} : \underline{u} \in \mathbb{U}^*\}^\perp \subset \mathcal{B}(\mathcal{H})_{sa}.$$

Vector space and positive cone:

$$V_0 := \mathcal{W}/K,$$

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Operator system lifts this to $V_0 \otimes \mathcal{B}(\mathbb{C}^n)_{sa}$:

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[Farenick/Paulsen, Math. Scand. 111:210-243, 2012]

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CP maps: $(\mathcal{D}_{\underline{u}} \otimes \text{id})C_n \subset C_n$ for all \underline{u} and n .

[Farenick/Paulsen, Math. Scand. 111:210-243, 2012]

But the \mathcal{D}_u^0 remember more than just being cp in the operator system. Indeed,

$$\mathcal{D}_u^0 \in \mathcal{P} := \{ \Lambda/K : \Lambda \text{ cp on } B(\mathcal{H}),$$

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$\mathcal{P} = \mathcal{P}(\mathcal{W}, K) \subset \mathcal{CP}(V_0)$, and in general

the inclusion is strict!

[Equality by Arveson's extension theorem for $K=0$ or $\mathcal{W} = B(\mathcal{H})_{sa}$]



6. Reconstructing the vector order?

Task: Find a suitable cone C for the g.u.r.
 $(V, \tau, \pi, \mathcal{D}_{\underline{u}})$, ideally a "nice" one...

Necessarily, $C_{\min} \subset C \subset C_{\max}$, with (recall)

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Can we choose C polyhedral or SDR?

Difficulty of course that C has to be preserved by the \mathcal{D}_u ; note that C_{\min} & C_{\max} satisfy this automatically.

6. Reconstructing the vector order?

Instructive special case: $C = C_{\min} = C_{\max}$,
ruling out a classical model if that is not
a polyhedral cone. [Cf. example, where this
happens with $C = \text{cone over a Bloch sphere.}$]

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ruling out a classical model if that is not
a polyhedral cone. [Cf. example, where this
happens with $C = \text{cone}$ over a Bloch sphere.
And the other example, where C is unique
and not SDR, in fact transcendental;
provides a process generated by a finite
GPT, but w/o quantum realisation.]