International Symposium on Quantum Information and Communication (ISQIC), 2025 CQUERE, TCG CREST

Kolkata, India





# Time: its strange aspects and how they reflect on quantum mechanics



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I'll give a quantum description of time based on conditional probability amplitudes

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I'll give a quantum description of time based on conditional probability amplitudes

... and an idea for a relativistic generalization: q spacetime

### WHAT is time?



### WHAT is time?

## In physics?





In physics?

### Time is what is measured by a clock





In physics?

### Time is what is measured by a clock

... but, what's a clock?!





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### ... or a "coordinate"



something that "measures" the distance between events

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something that "measures" the distance between events

the two **main** meanings of "time" in physics

# other meanings?!

#### Table 2.1: Times.

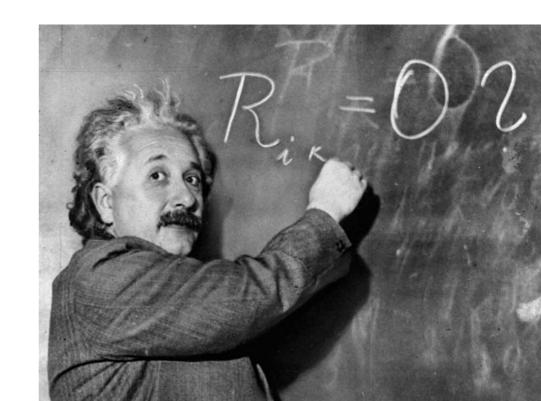
Time notion	Property	Example	Form
Natural language time	memory	brain	?
Time-with-a-present	present	biology	R
Thermodynamical time	direction	thermodynamics	A
Newtonian time	unique	newtonian mechanics	M
Special relativistic time	external	special relativity	$M^3$
Cosmological time	spatially global	cosmological time	m
Proper time	temporally global	world line proper time	$m^{\infty}$
Clock time	metric	clocks in GR	c
Parameter time	one dimensional	coordinate time	$L^{\infty}$
No-time	none	quantum gravity	none

[Rovelli, "quantum gravity"]

Oclock time Laproper time VZMXA Oquantum time a parameter O coordinate time X @ time of arrival @Anthropic time B Waw time >NO time and menory O "time" time DLeibni

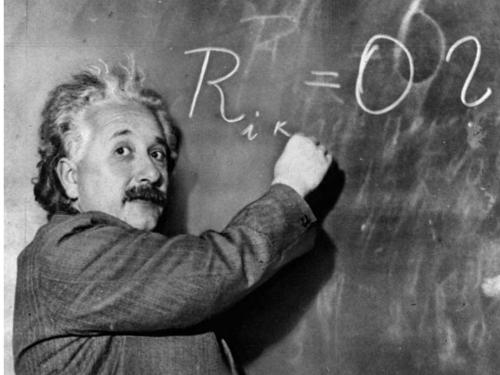


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#### rather: past-present-future have different degrees of existence (whatever that means)

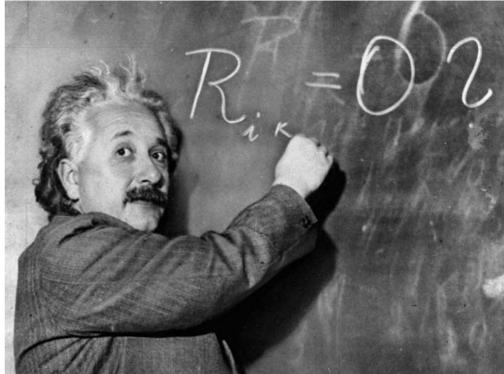


special relativity

NO!

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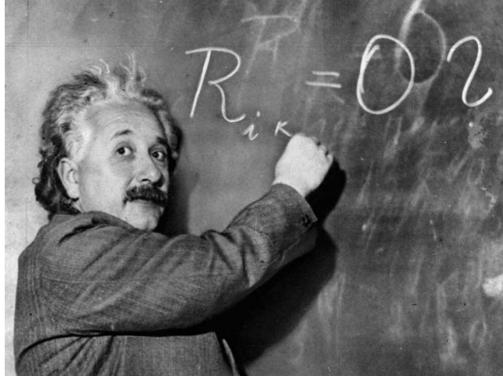
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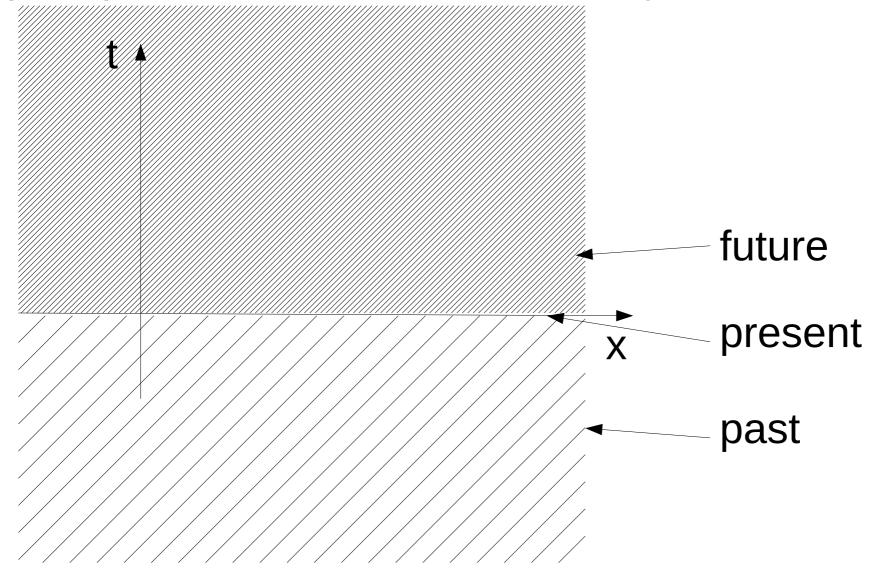


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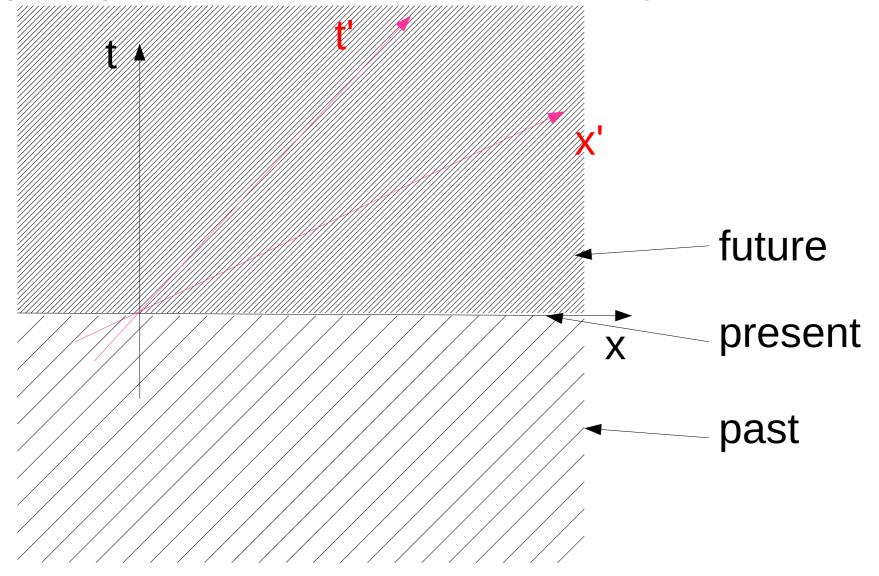
"NO"?!? why?



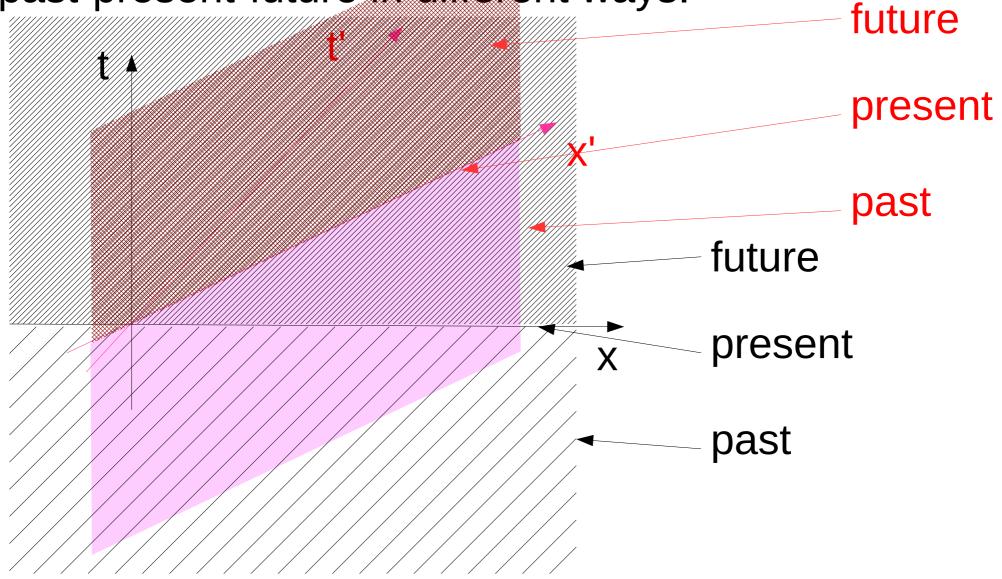
# Observers in relative motion divide spacetime in past-present-future in different ways.



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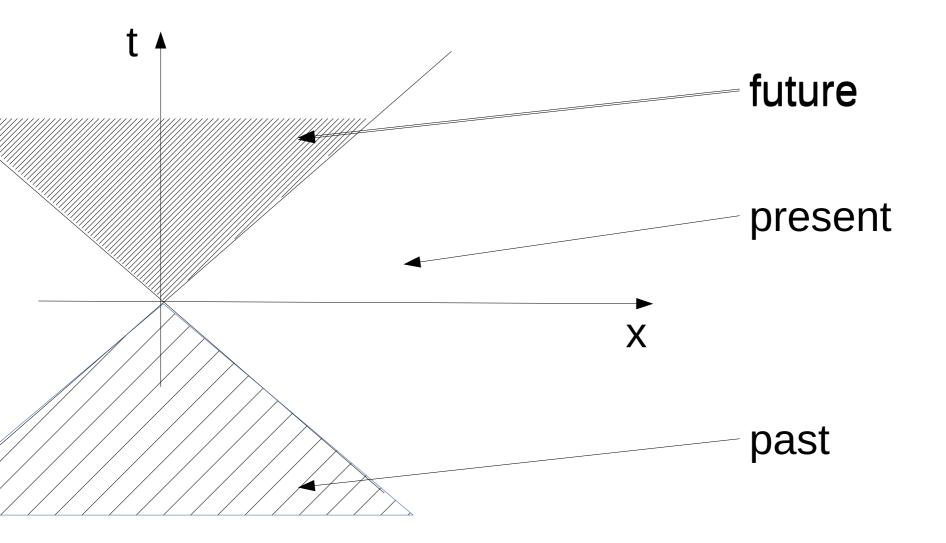


Observers in relative motion divide spacetime in past-present-future in different ways.

### The present is **relative to the observer** So, whose present should be "real"!?

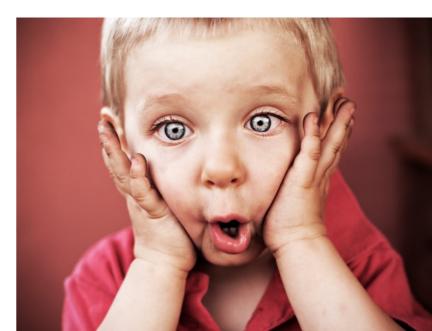
present

# Light cones **are** invariant: can we use them to define past-future-present?



# Light cones **are** invariant: can we use them to define past-future-present? **NO!!!** future present Х past

Now even observers in the same reference disagree on what is "real"



# Some people (e.g. Lee Smolin) disagree, but they usually reject Galileian relativity



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(not incompatible with observations, but incompatible with textbook relativity).



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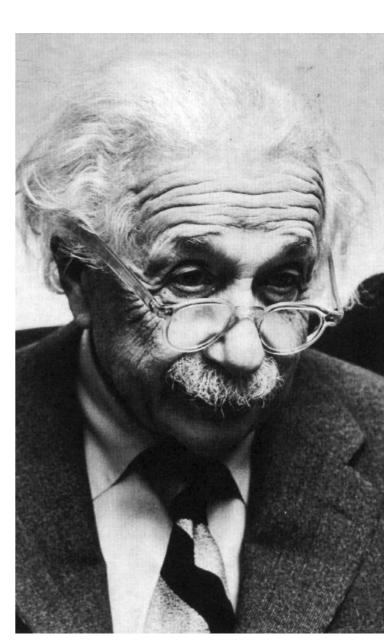
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Technically: Presentism vs Eternalism



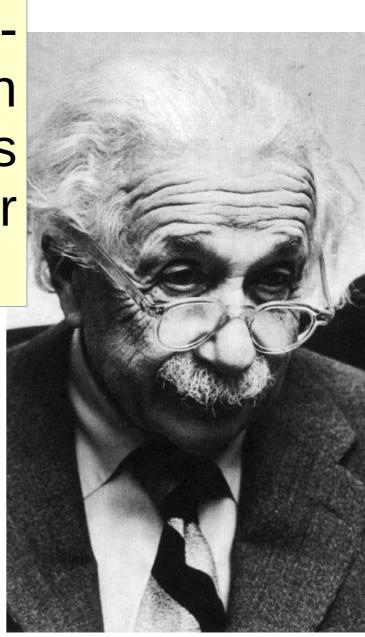
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"For us convinced physicists the distinction between past, present, and future is only an illusion, however persistent."

Albert Einstein, 21 May 1955

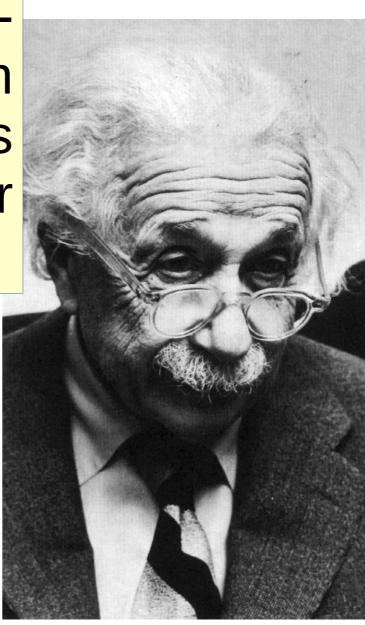


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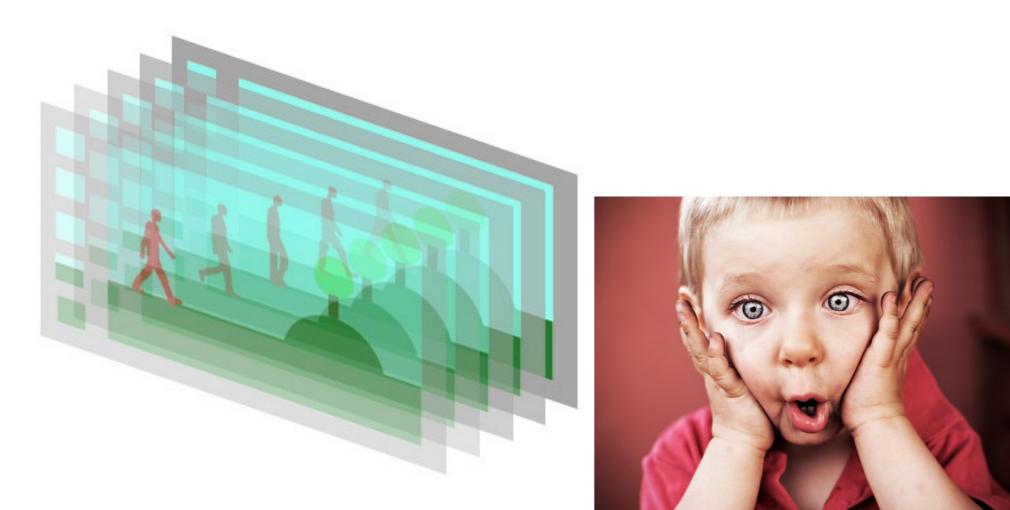
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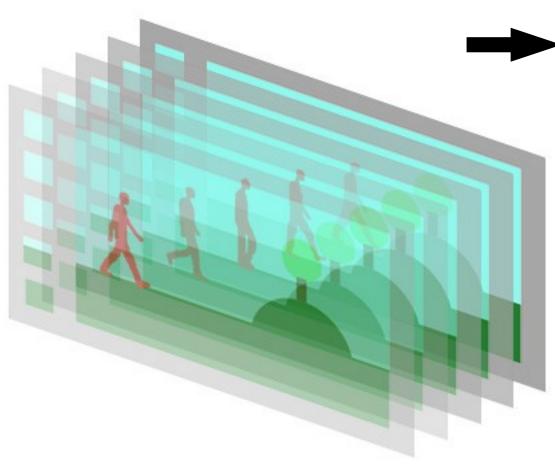
In a letter to the widow of his dear friend Michele Besso: trying to console her (or himself?) with special relativity.



# Consequence: Block universe



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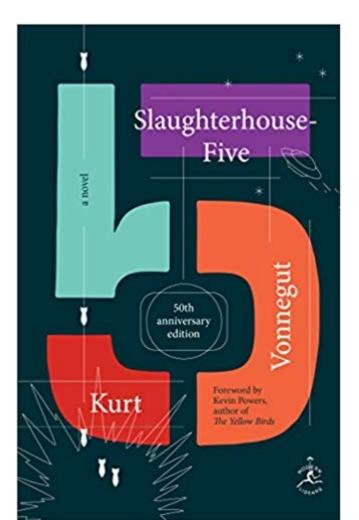
we'd like a quantum description of time that contains the BU



# Our intuition fails badly...



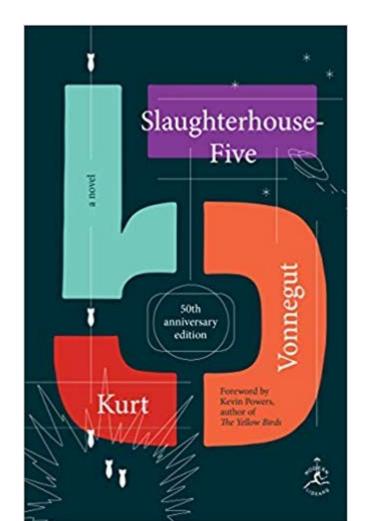
"Story of Your Life" by Ted Chiang.



# Our intuition fails badly...

our perception of time is incompatible with relativity:

- we perceive time locally (only the present)
- we perceive space globally (we don't perceive only our own location)



#### ...join GR and QM?!?

Canonical quantization of GR

## Wheeler-De Witt equation:

# $\hat{H}|\Psi\rangle = 0$



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Canonical quantization of GR

Wheeler-De Witt equation:

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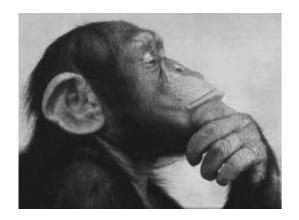
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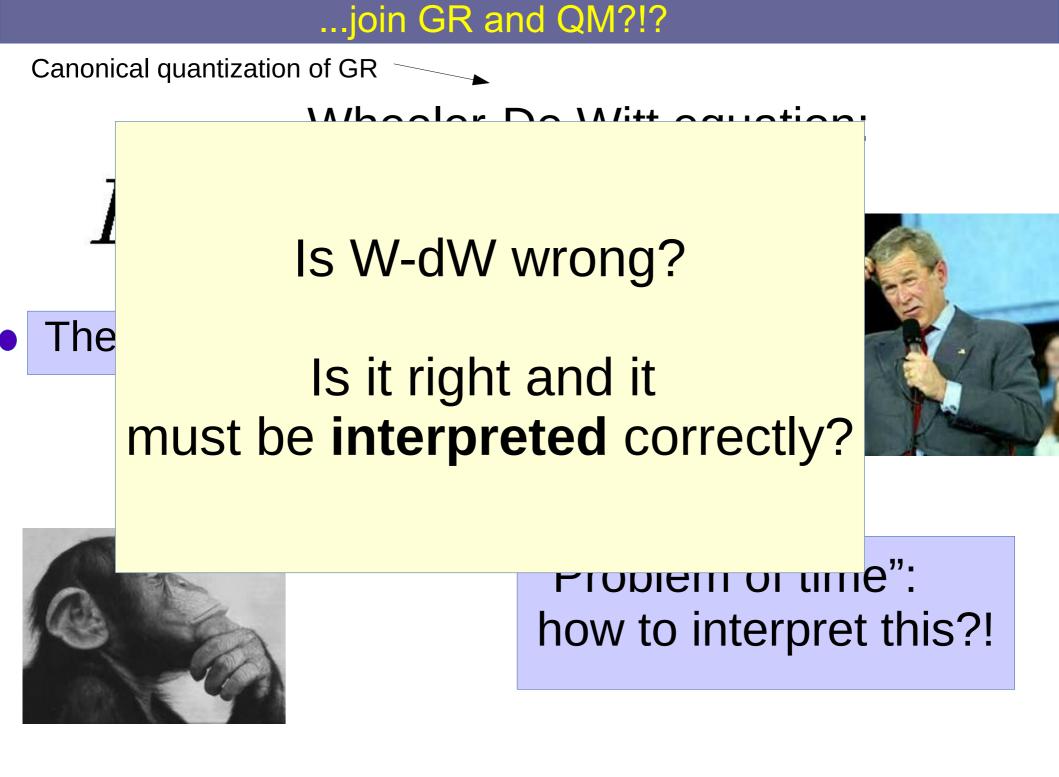
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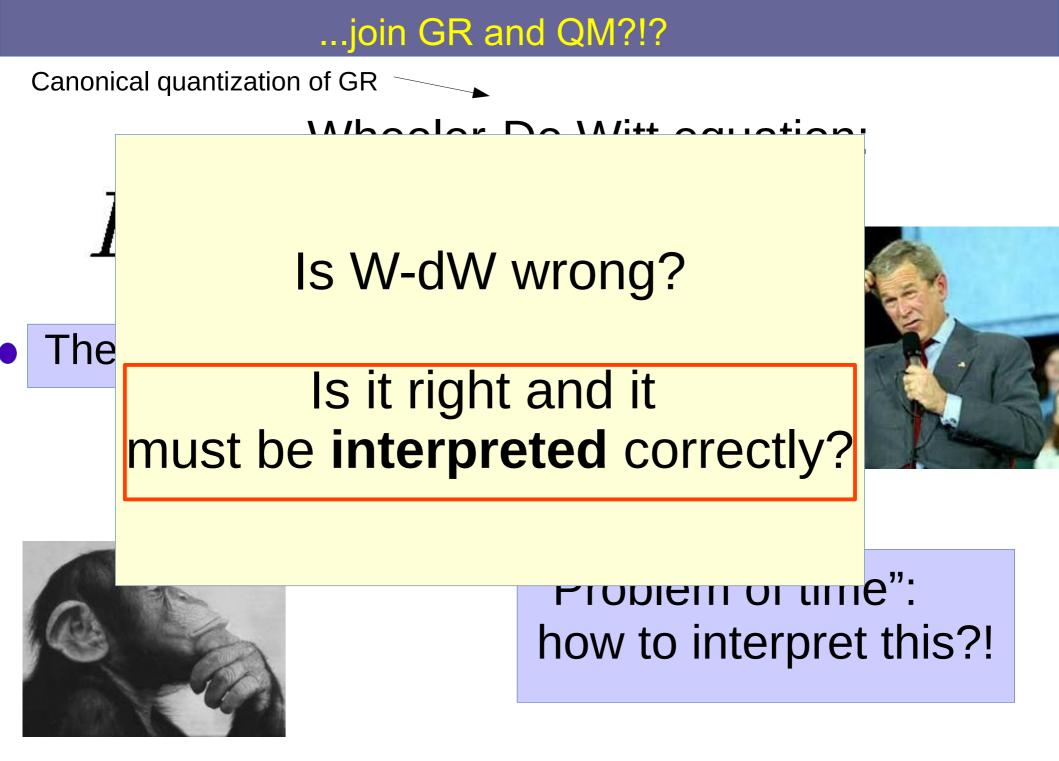




### ...but!!!

### "Problem of time": how to interpret this?!







# Take a step back... Time in non-relativistic QM



# Time in quantum mechanics:



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a classical parameter in the Schroedinger eq.

$$i\hbar\frac{\partial}{\partial t}|\psi(t)\rangle = \hat{H}|\psi(t)\rangle$$



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$$i\hbar\frac{\partial}{\partial t}|\psi(t)
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it indicates what is shown on the **clock** on the lab wall.



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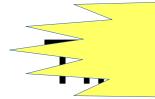
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# BUT... classical systems don't exist in a consistent theory of

(they're just a limiting situation)



Inconsistency in the formulation of QM (inconsistency in one of its postulates)

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#### **BUT** --- classical systems don't exist in a consistent theory of quantum mechanics

in a consistent theory of quantum mechanics (they're just a limiting situation)

**Quantum Time** 

# define: Time is "what is shown on a clock"

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## hen use a quantum system as a clock

### e.g. a quantum particle on a line (or any other quantum system)

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 $\mathcal{H} \equiv \mathcal{L}^2(\mathbb{R})$  eigenbasis  $\{|x\rangle\}$ 

#### **Time and entanglement**



# Time arises as **correlations** between the system and the clock



Page and Wootters [PRD **27**,2885 (1983)] Aharonov and Kaufherr [PRD **30**, 368 (1984)]

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This means that for physical states the system Hamiltonian is the generator of *clock* time translations

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• to the energy being  $\omega$ :

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I = T

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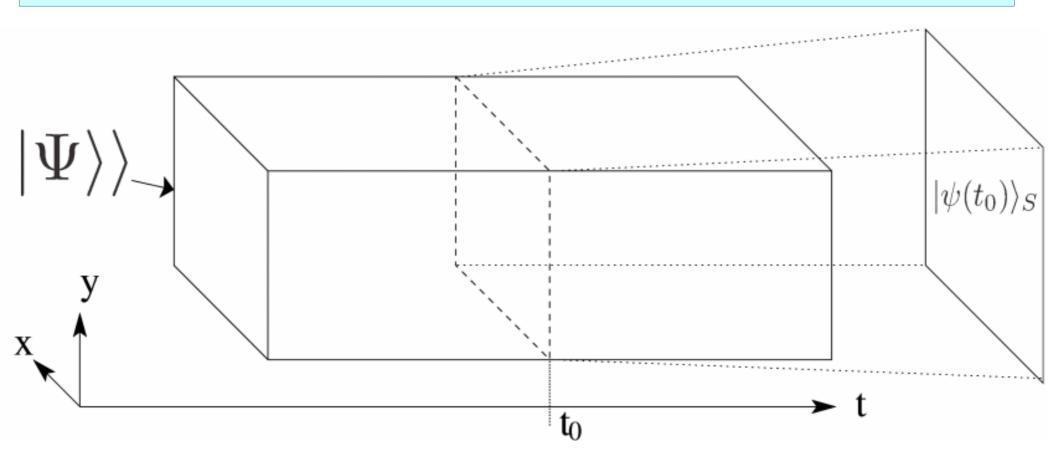
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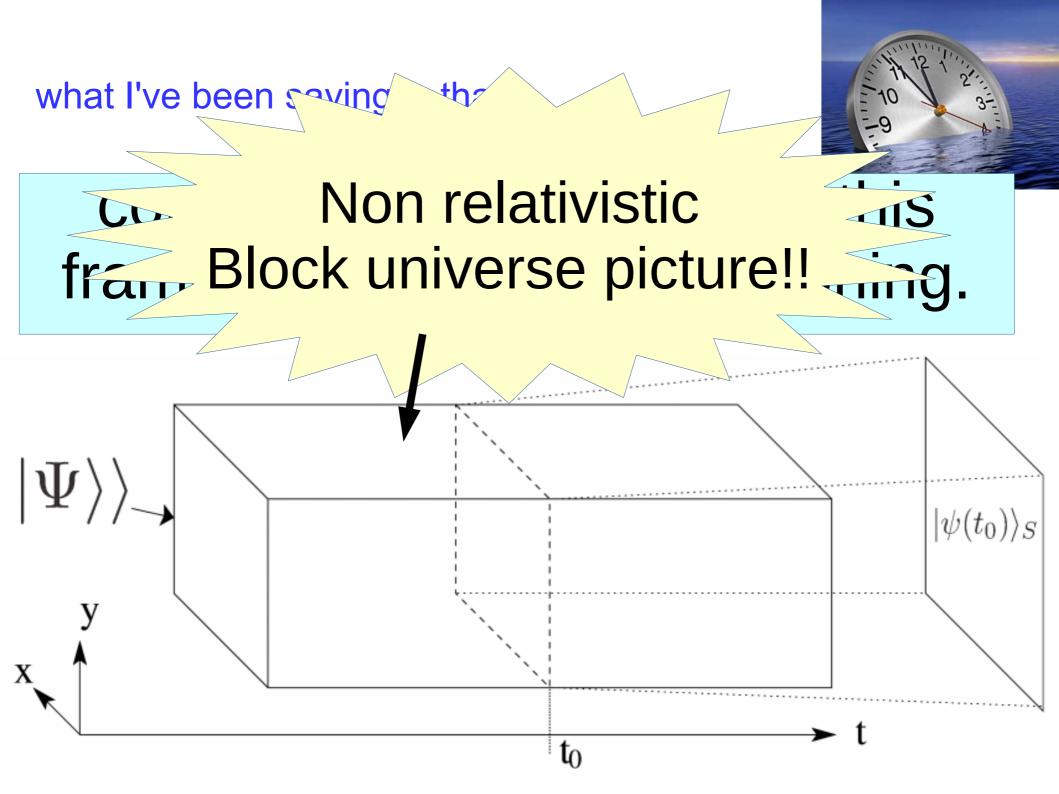
"momentum" representation=time indep. Schr eq.

what I've been saying is that



# conventional qm arises in this framework through conditioning.





# conditioning?

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are of the form:

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## conditioning?

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Remember this! Relation QM-Relativity

which mean system at me

are of th

765

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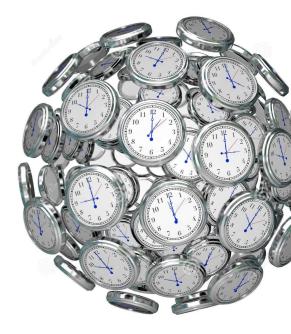
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 $dt t |t\rangle$ 

(it's a **continuous** quantum degree of freedom with the choice  $\mathcal{H} \equiv \mathcal{L}^2(\mathbb{R})$ )  $\hat{\mathcal{I}}$ Other choices are possible!! **Physical interpretation** 

## The time Hilbert space is the Hilbert space of the clock that **defines** time

remember: "time is what is measured by a clock"!



**Physical interpretation** 

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**Physical interpretation** 

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here: we used a Hilbert space for a particle on a line, appropriate for a continuous time that goes from  $-\infty$  to  $+\infty$ 

other choices are possible..

if the clock has finite energy, time is cyclic: e.g. a spin (appropriate for certain closed cosmologies) Up to now: the time Hilbert space is the Hilbert space of the clock that **defines** time

## BUT, a physical interpretation of the time Hilbert space is **un-necessary**



Up to now: the time Hilbert space is the Hilbert space of the clock that **defines** time

BUT, a physical interpretation of the time Hilbert space is **un-necessary** 

alternative:



It can be seen as an **abstract** purification space

Is entanglement important? Could we do with classical correlations?

$$\begin{split} |\Psi\rangle\rangle &= \int dt \; |t\rangle_T \otimes |\psi(t)\rangle_S \\ &= \int d\mu(\omega) \; |\omega\rangle_T \otimes |\psi(\omega)\rangle_S \; , \end{split}$$



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**NO!**



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### Instead of bipartite entanglement

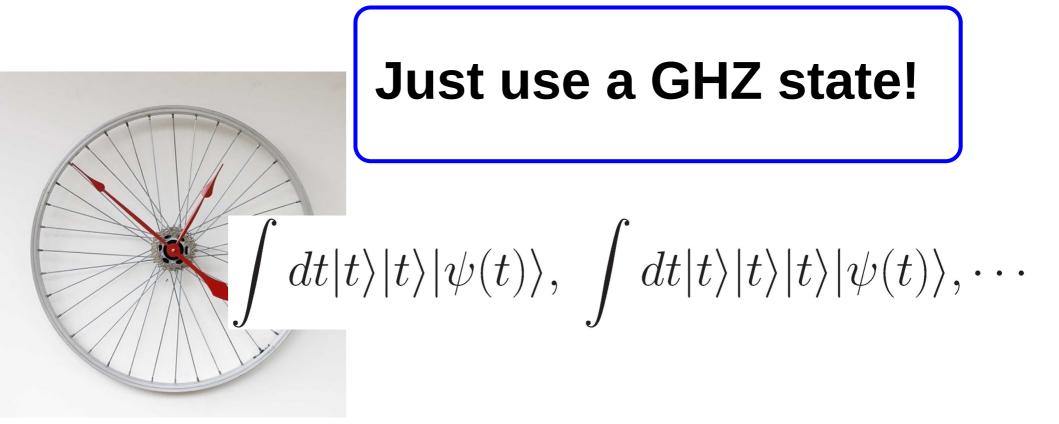
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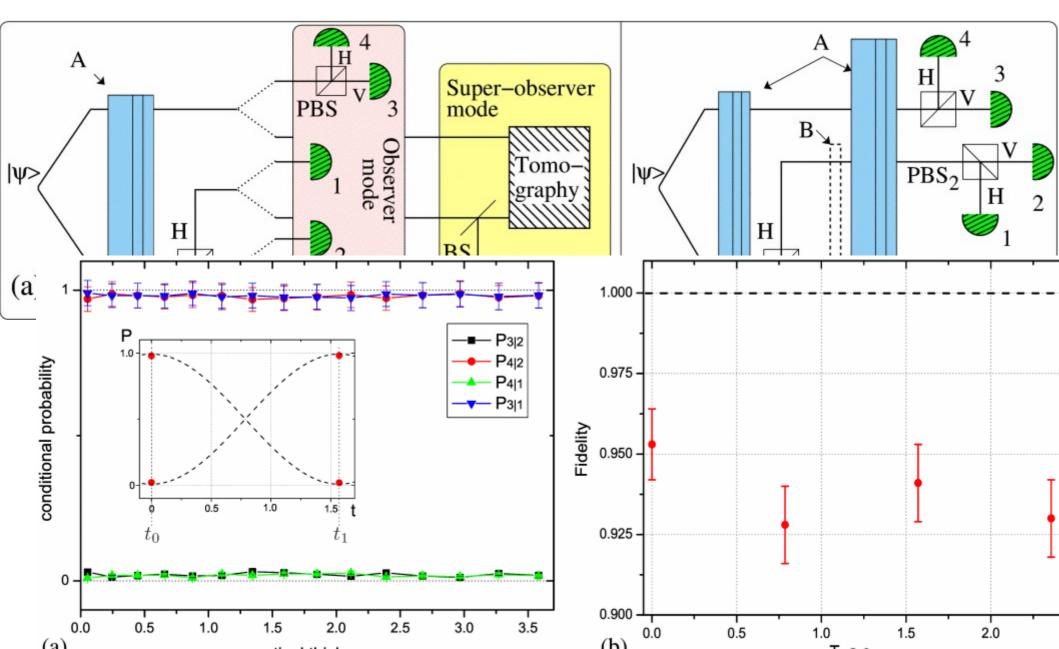
### Just use a GHZ state!

### Instead of bipartite entanglement

$$|\Psi\rangle\rangle = \int dt \; |t\rangle_T \otimes |\psi(t)\rangle_S$$



## Experimental illustration (collaboration with the INRIM group)





These ideas were basically abandoned in the 80s: because of objections (Kuchar, Unruh, etc.)

What Ever

We removed these objections

... and also perfected the model (e.g. role of entanglement, momentum representation)

# Criticisms to time quantizations

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... but wait!! In our case we have  $\begin{bmatrix} \hat{T} & \hat{O} \end{bmatrix} = i\hbar \rightarrow \lambda(\hat{O}) \subset (-\infty)$ 

$$[T,\Omega] = i\hbar \Rightarrow \lambda(\Omega) \in (-\infty, +\infty)$$

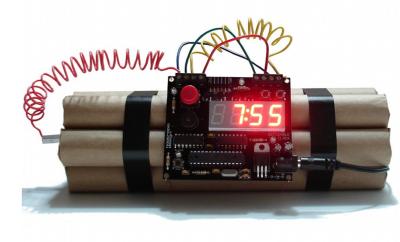
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NOT the system Hamiltonian  $H_S$  !!!

can be anything In other words, the **Pauli argument fails** in our case because the energy-time connection is not enforced dynamically as

$$[\hat{T}, \hat{H}_S] = i\hbar$$

but as a constraint on the physical states through a WdW eq:  $\hat{\mathbb{J}}|\Psi
angle
angle=0$ 

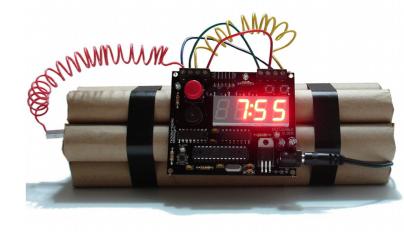


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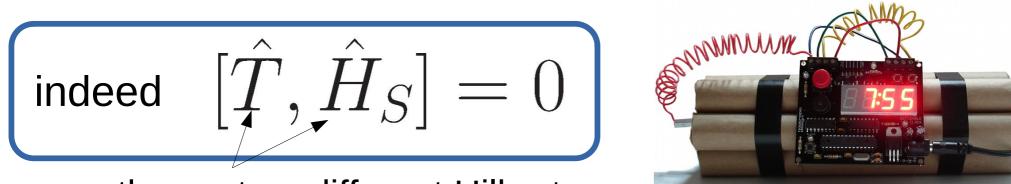
indeed 
$$[\hat{T},\hat{H}_S]=0$$



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they act on different Hilbert spaces

#### The Peres argument

Peres argument: "if energy generates time translations and momentum generates position translations, then the Hamiltonian and the momentum operator should commute always"

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•in conventional qm, time is not a dynamical variable  $\Rightarrow$  no problem.



• in our case, time is a dynamical variable, but its translations are NOT generated by  $\hat{H}_S$  (but by  $\hat{\Omega}$ )

#### The Kuchar argument against PaW

Kuchar: "measurements of a system at two times will give the wrong statistics: the first measurement "collapses" the time d.o.f. and the system remains stuck"



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## Kuchar's objection killed PaW's argument

Kuchar: "measurements of a system at two times will give the wrong statistics: the first measurement "collapses" the time d.o.f. and the system remains stuck"

$$\begin{split} |\Psi\rangle\rangle &= \int dt \; |t\rangle_T \otimes |\psi(t)\rangle_S \\ & \int \text{time } t \\ |\psi(t)\rangle \end{split}$$



after a measurement of time, the state collapses to  $|\psi(t)\rangle$  : successive measurements give wrong statistics: no more evolution

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## a careful formalization of **what a two-time measurement is** solves the problem!

Kuchar: "measurements of a system at two times will give the wrong statistics: the first measurement "collapses" the time d.o.f. and the system remains stuck"

a careful formalization of **what a two-time measurement is** solves the problem!

The second measurement is a joint measurement on the system and on the d.o.f. that stored the outcome of the first.



In formulas (using von Neumann's prescription for a measurement):

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$$|\psi(t_0)\rangle_S|\mathbf{r}\rangle_m \xrightarrow{U} |\psi'\rangle_{Sm} \equiv \sum_a \psi_a |a\rangle_S|\mathbf{a}\rangle_m$$

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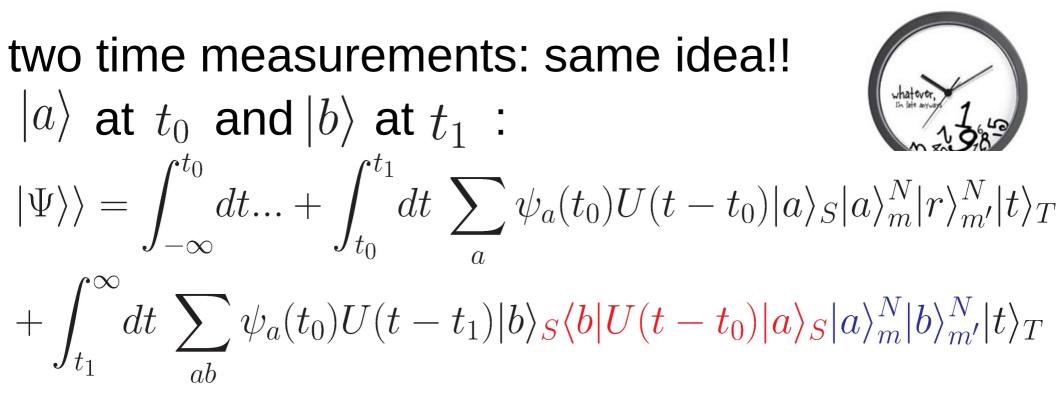
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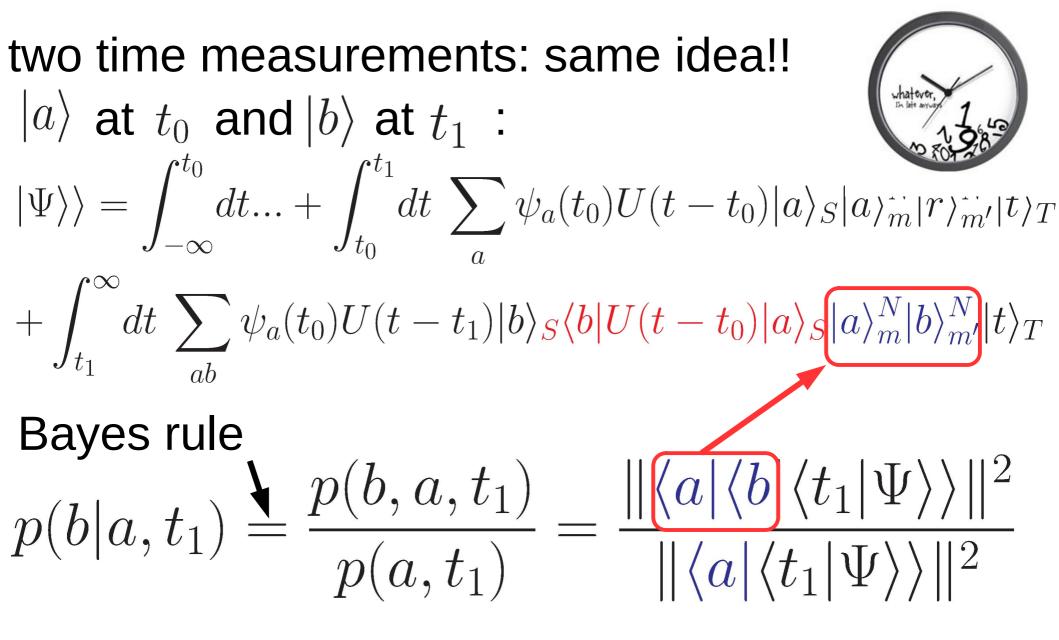
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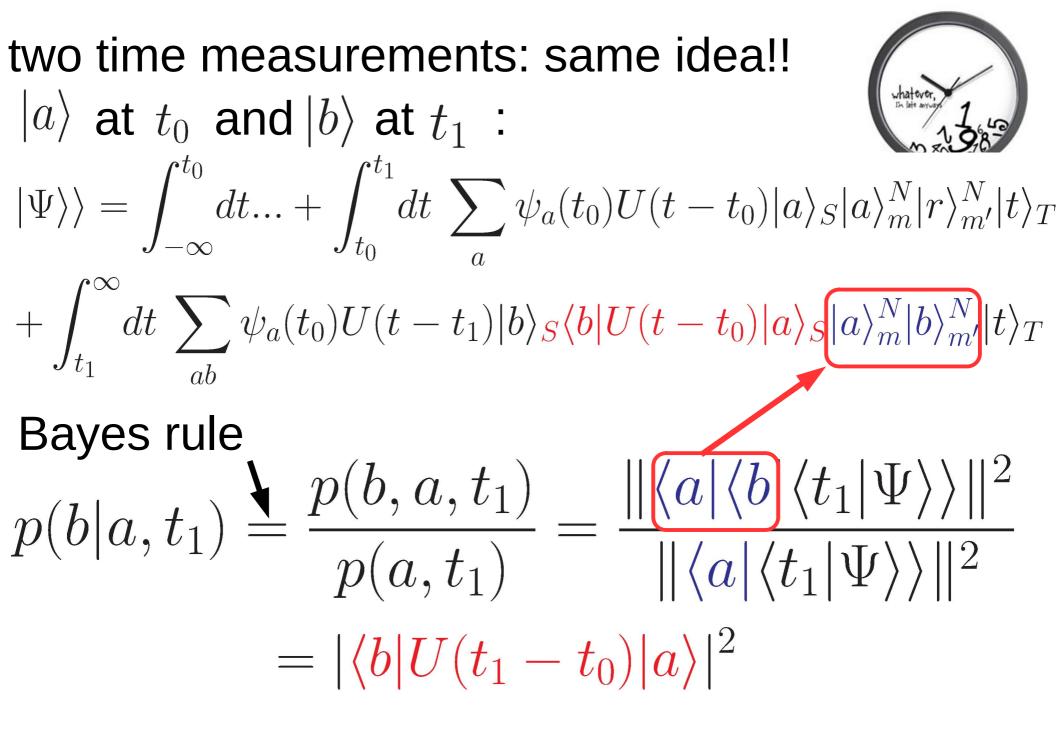
$$\Rightarrow p(a|t_0) = |\langle a|\psi(t_0)\rangle|^2 \equiv ||_{m} \langle a|_{T} \langle t_0|\Psi\rangle\rangle|^2$$
$$= |\psi_a(t_0)|^2 \quad \text{(Born's rule)}$$

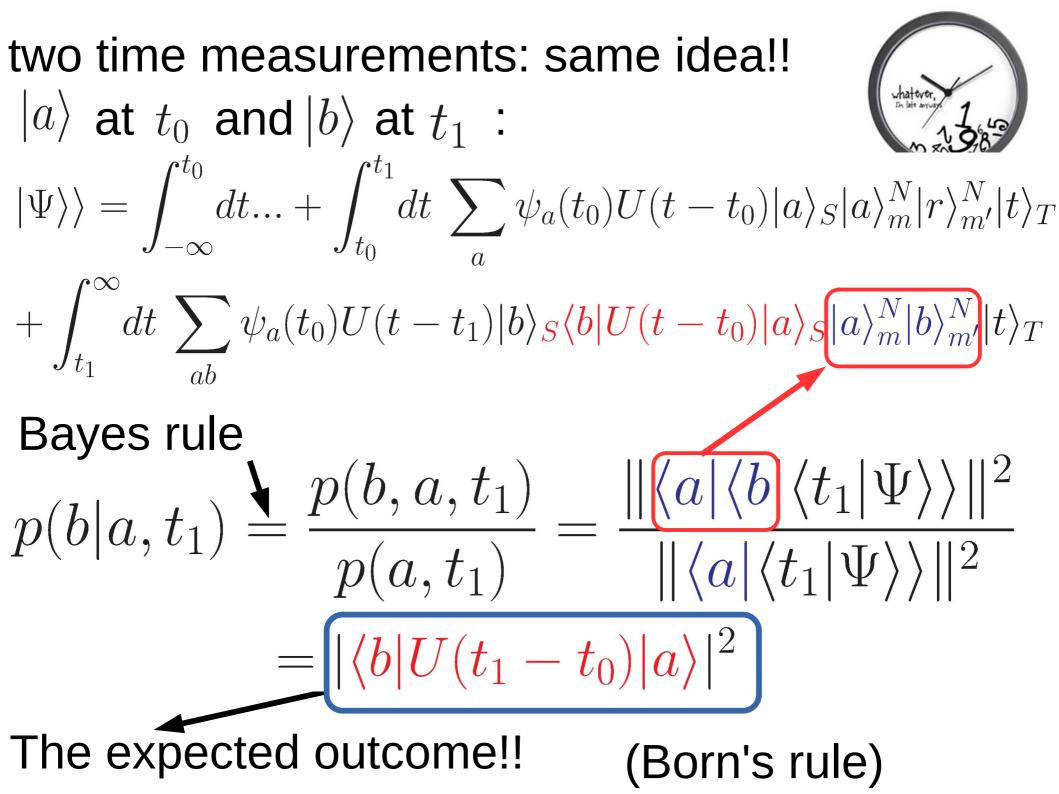
### two time measurements: same idea!!

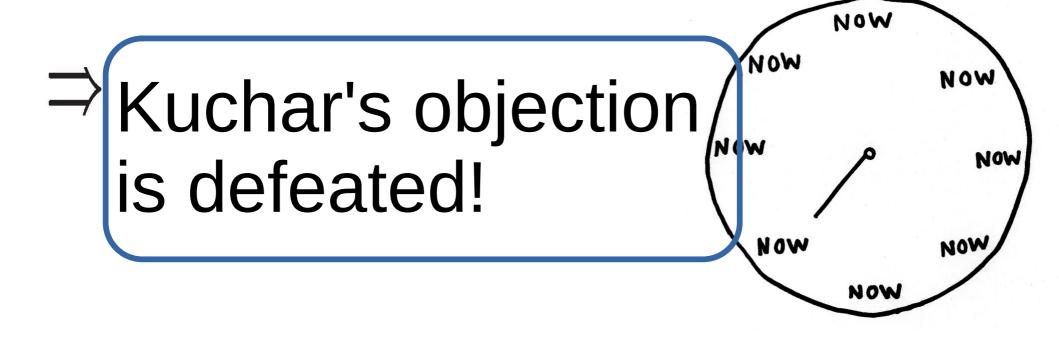




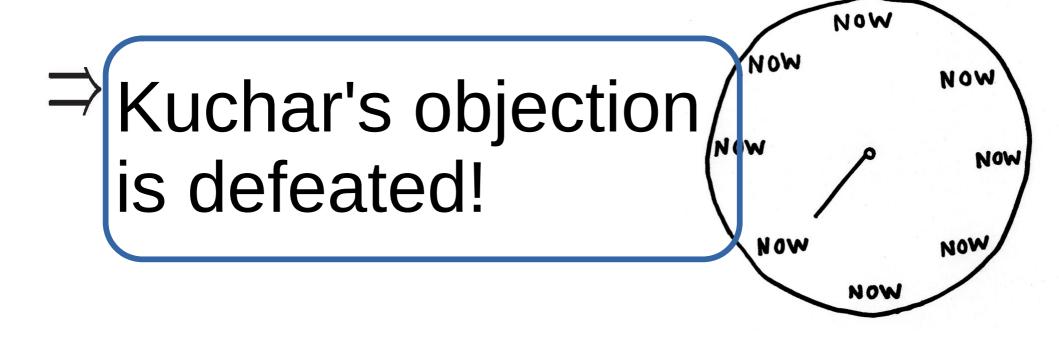












this argument can be extended to POVMS, propagators, etc...





What are the hypotheses for this argument? Use von Neumann's quantum mechanics! (Born's rule and all that)



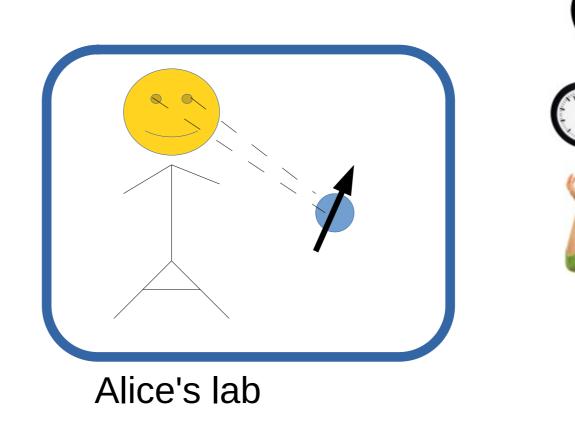
## Use von Neumann's quantum mechanics! (Born's rule and all that)

While we do admit that a unitary description of a measurement apparatus must exist, we still work in the conventional quantum framework.



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**Bob's point** of view Alice's lab

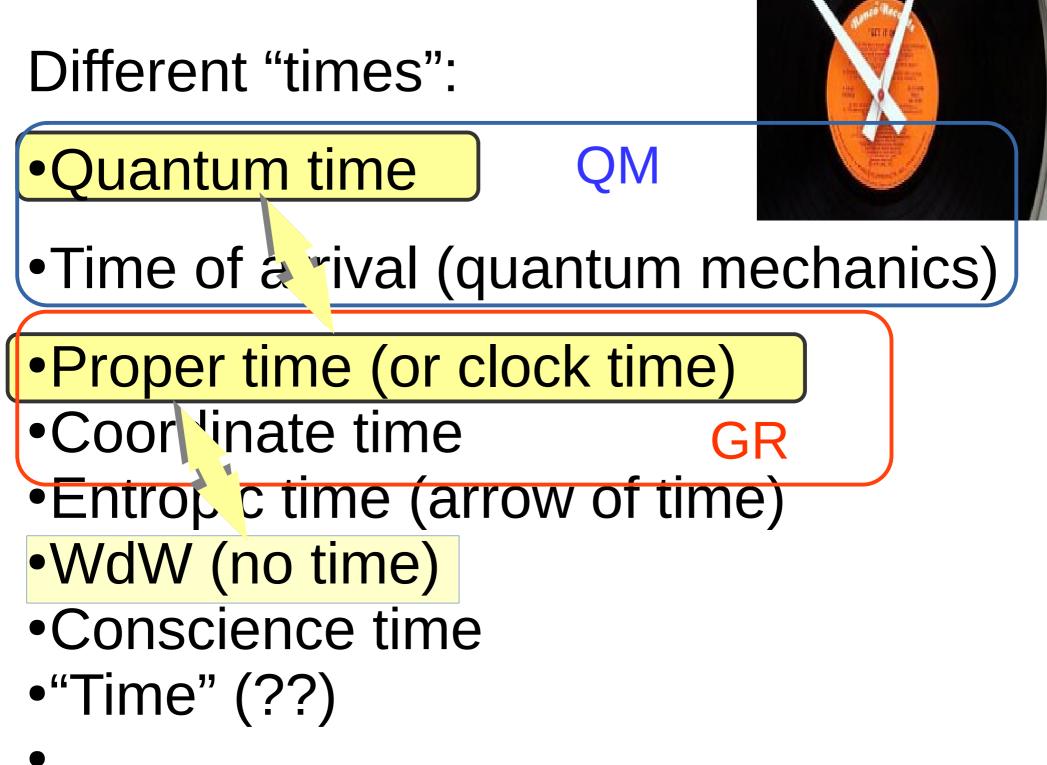
- Different "times":
- Quantum time



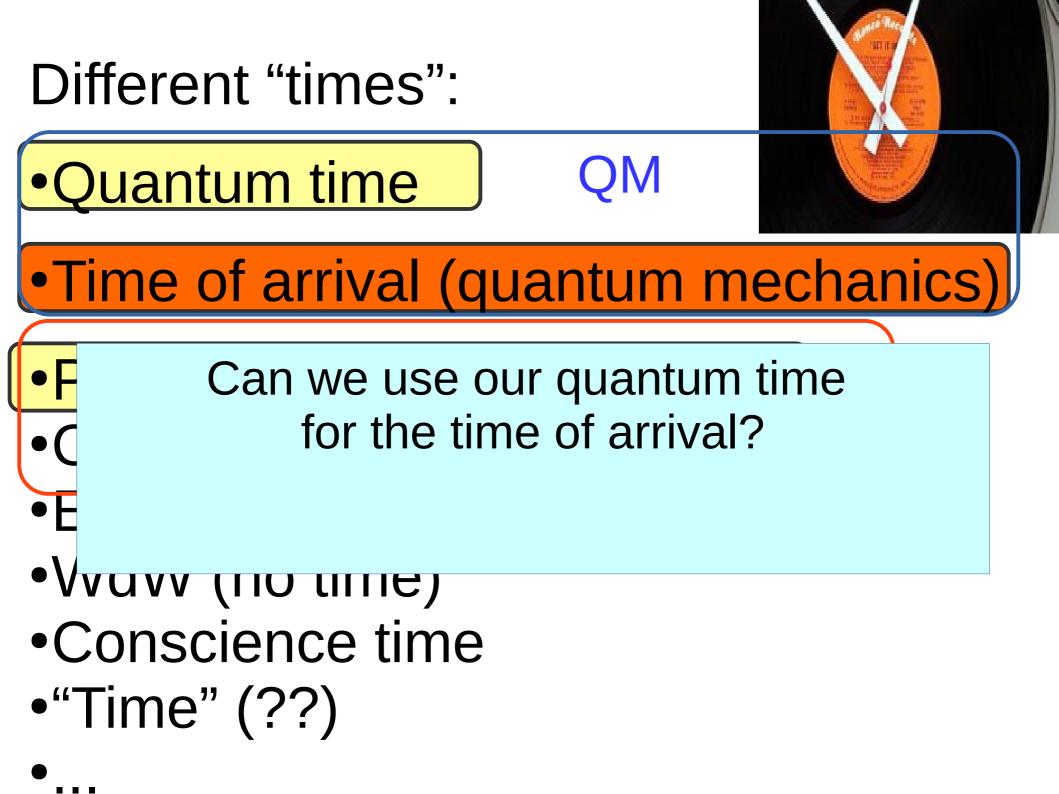
- •Time of arrival (quantum mechanics)
- Proper time (or clock time)
- Coordinate time
- •Entropic time (arrow of time)
- •WdW (no time)
- Conscience time
- •"Time" (??)
- •

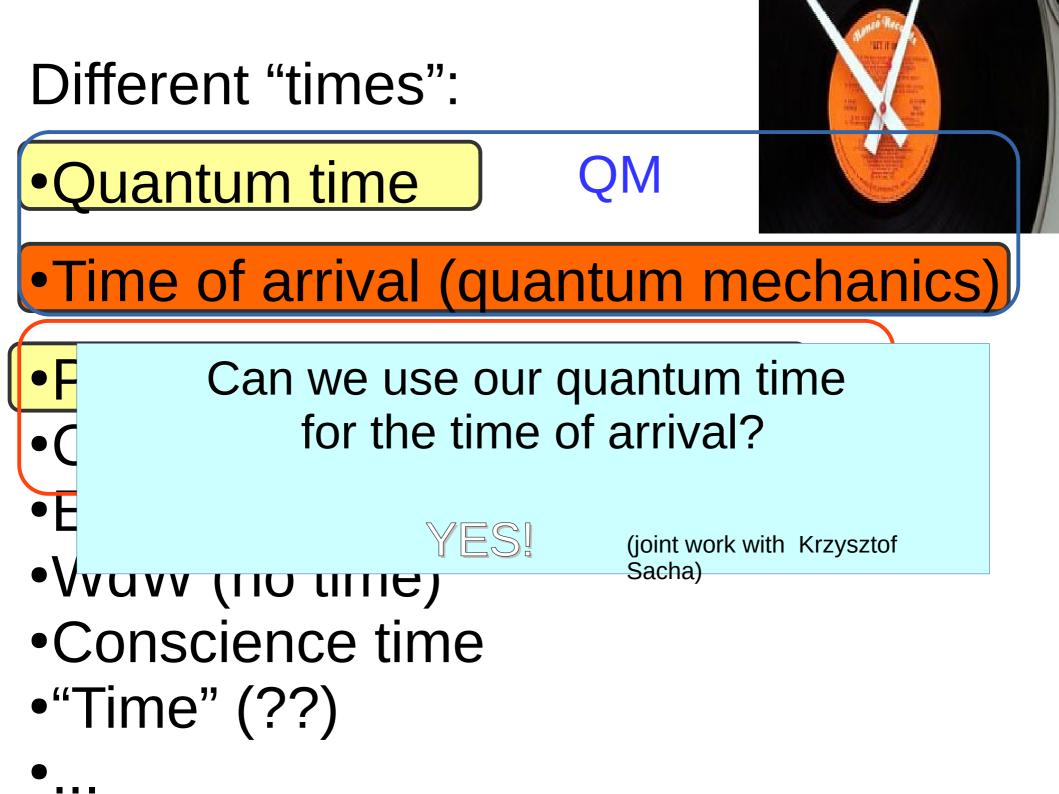


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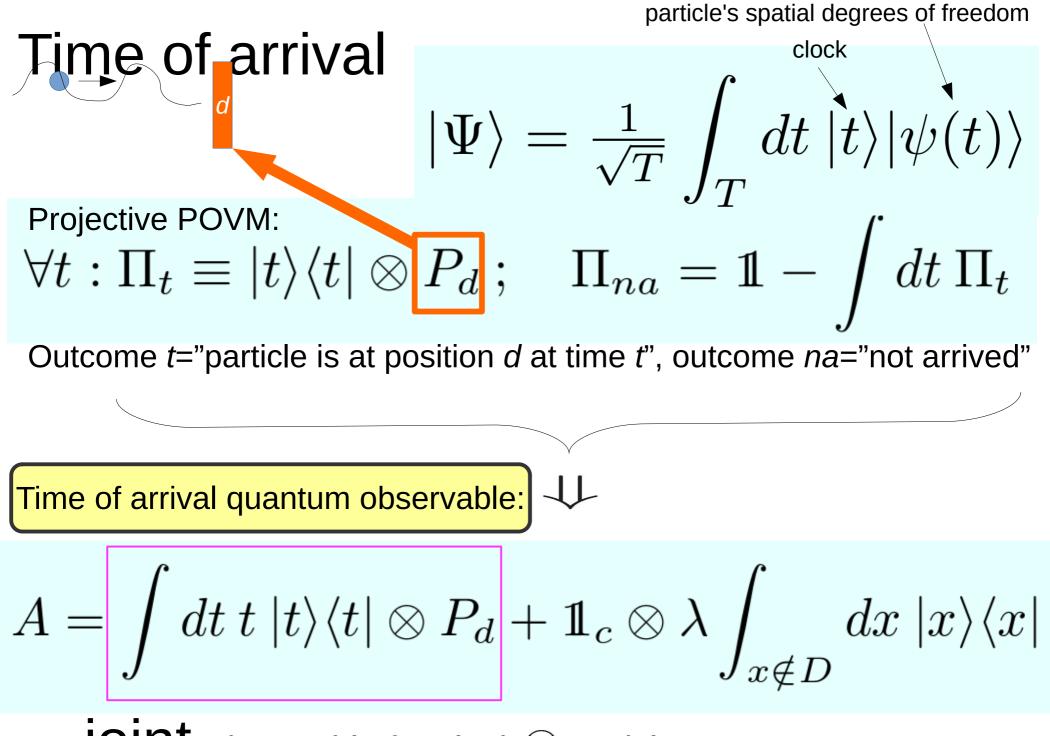


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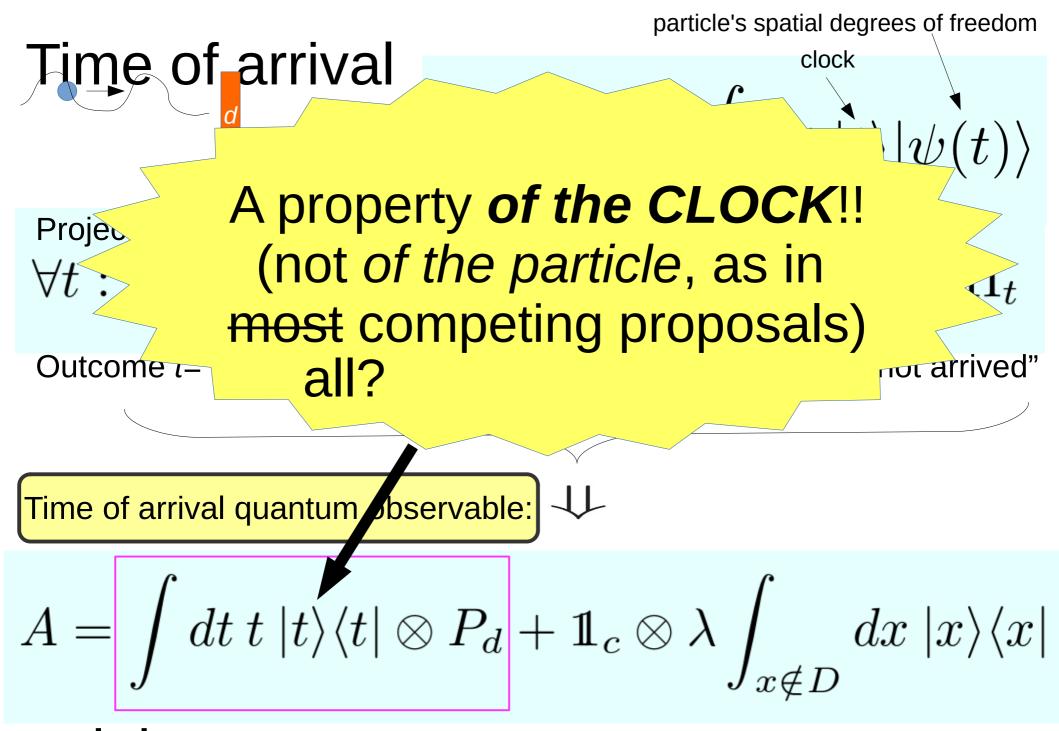








A **JOINT** observable for clock  $\otimes$  particle



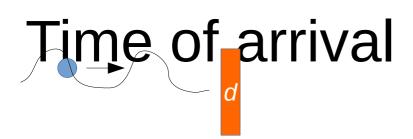
A **JOINT** observable for clock  $\otimes$  particle

particle's spatial degrees of freedom Time of arrival clock  $|\Psi\rangle = \frac{1}{\sqrt{T}} \int_{T} dt \, |t\rangle |\psi(t)\rangle$ **Projective POVM:**  $\forall t : \Pi_t \equiv |t\rangle \langle t| \otimes P_d; \quad \Pi_{na} = \mathbb{1} - \int dt \, \Pi_t$ Outcome *t*="particle is at position *d* at time *t*", outcome *na*="not arrived"  $(t, x = d) = \operatorname{Tr}[|\Psi\rangle\langle\Psi|\Pi_t] = \frac{1}{T}|\psi(d|t)|^2$ , Born's rule with  $\psi(x|t) \equiv \langle x|\psi(t)\rangle$  $p(t|x=d) = \frac{p(t,x)}{2}$  $\left|\frac{x=d}{x}\right| = \frac{|\psi(d|t)|^2}{\int_T dt \, |\psi(d|t)|^2} ,$ yes rule Time of arrival prob. distribution











• take the projector for the particle at *d* and for the clock at *t*.

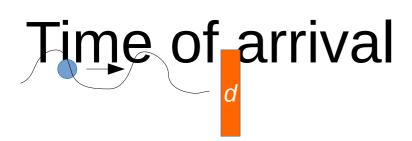






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- build a joint observable from this







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- take the projector for the particle at *d* and for the clock at *t*.
- build a joint observable from this
- from the joint probability of clock+particle, get the time of arrival prob through the Bayes rule.

# Only "time of arrival"?



# Only "time of arrival"? NO!

Extensions to other time measurements:

# a generic time measurement is

# "At what time did the event E happen?"

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t<sub>E</sub>

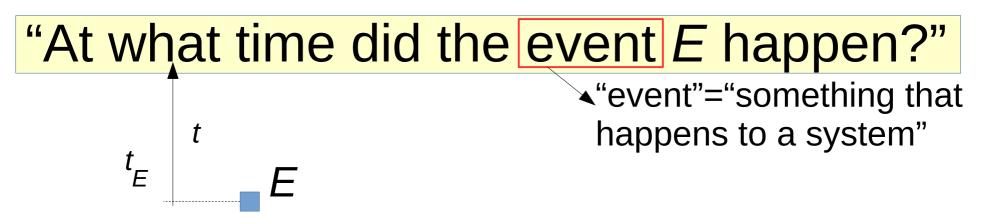
# "At what time did the event E happen?"

"event"="something that happens to a system"

# Only "time of arrival"? NO!

Extensions to other time measurements:

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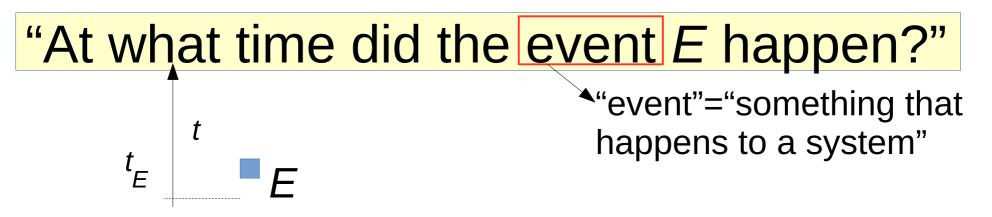


Use the same trick: a joint projector on the clock and on the system (the projector on the system referring to the event *E*)

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e.g. "at what time is the spin up?" The projector is  $|\uparrow\rangle\langle\uparrow|$ 

All usual manipulations

for observables can be done:

- Expectation values
- Probability distributions
- •Eigenstates, eigenvalues, etc.



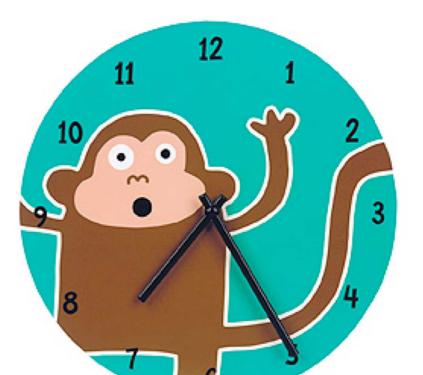
Advantages with respect to previous proposals:



- Describe situations that prev prop could not (multiple pass, stationary particle, etc.)
- •Extension to arbitrary events
- Possibility of describing multiple clocks
- Testable differences: experiment!

Up to now: nonrelativistic QM

Can we use similar ideas for relativistic QM?





### •GR — • events



### quantum systems

Inifinitely extended in time (finite or infinite in space)

### events

•GR

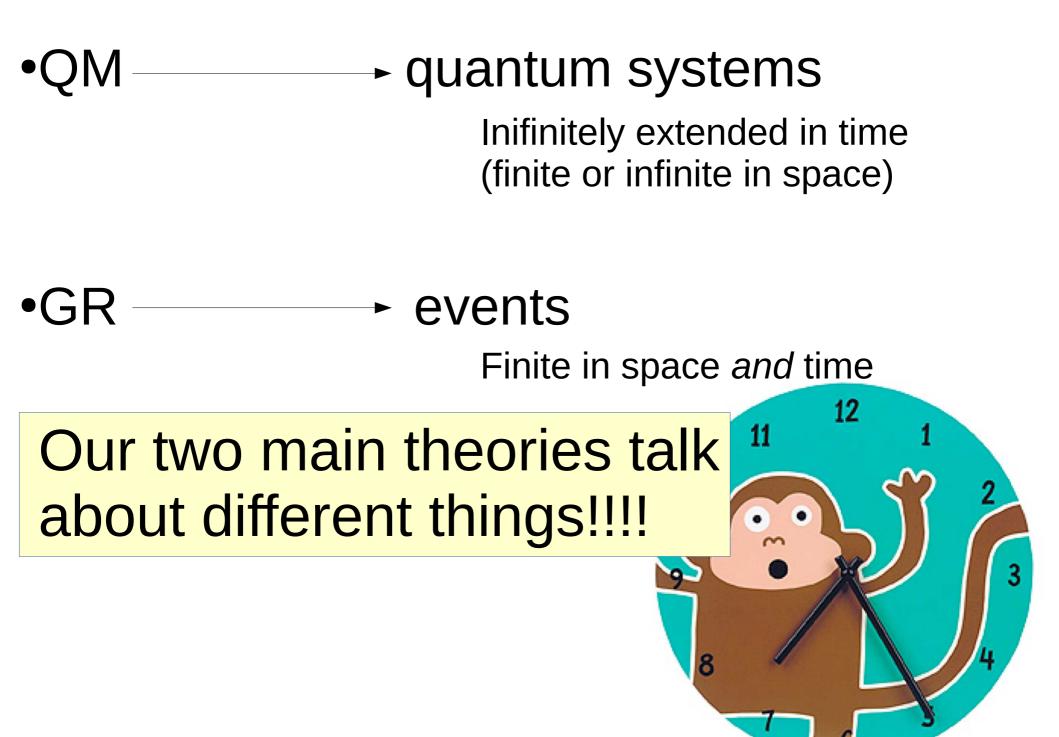
Finite in space and time

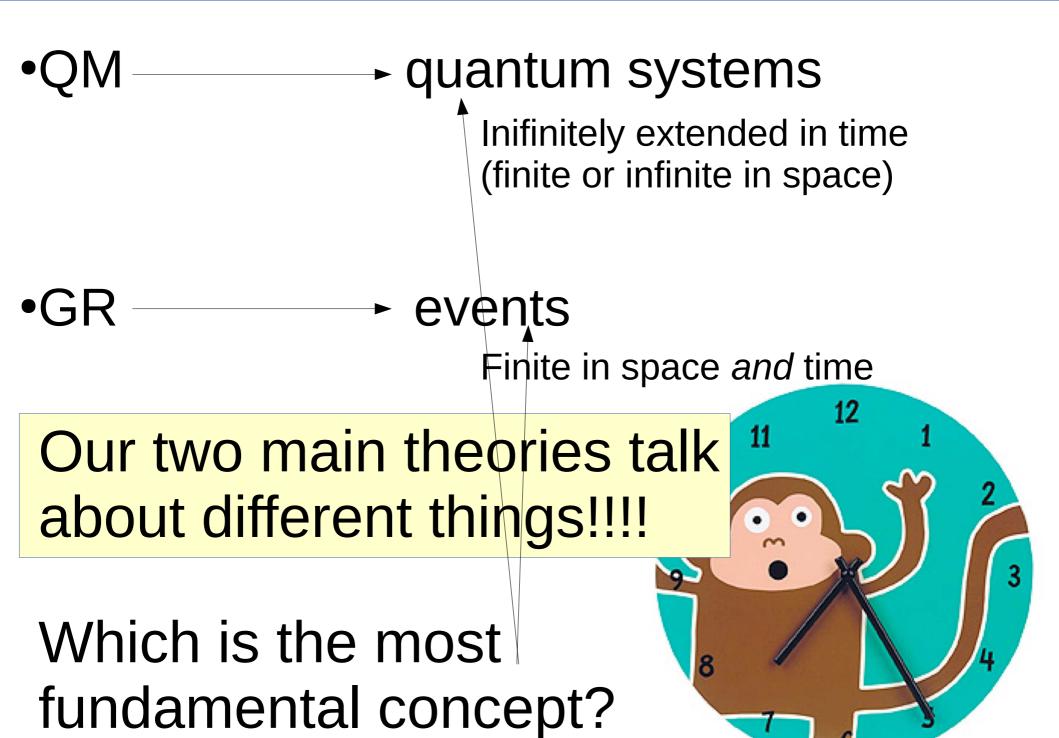
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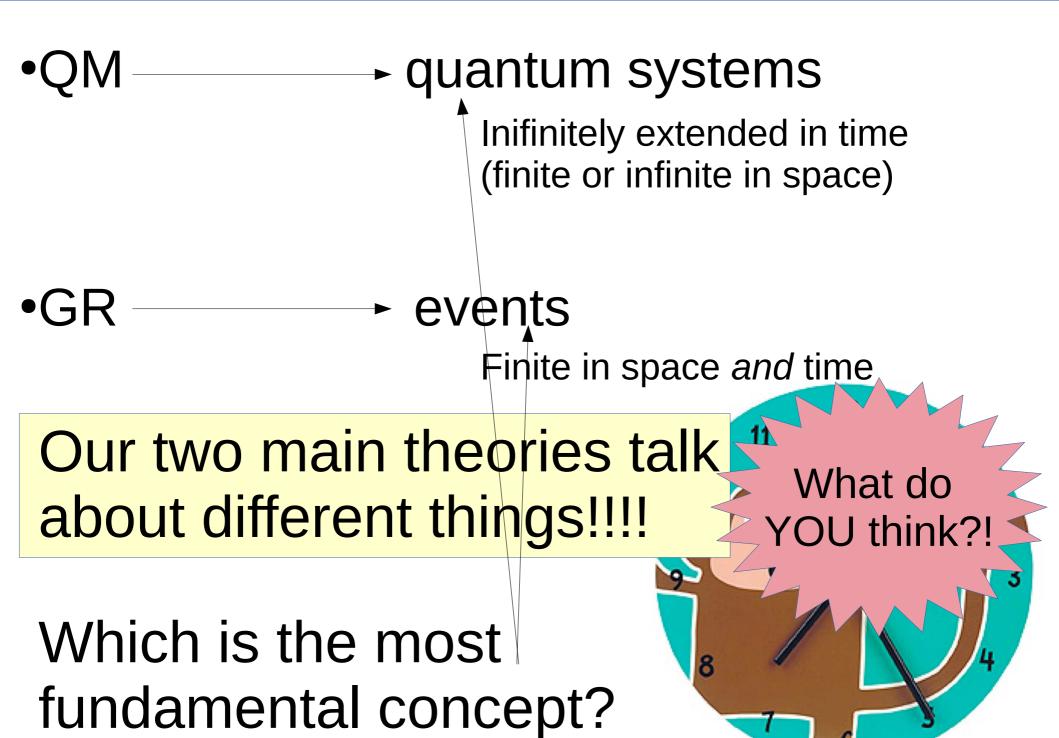
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### General relativistic theory of QM

Systems are more fundamental: GR is made of quantum fields

### •Quantum theory of GR

Events are more fundamental:



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Events are more fundamental:

Quantum system=succession of events



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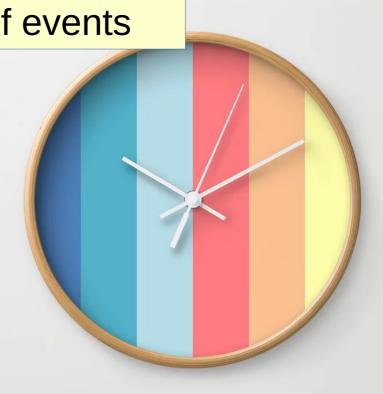
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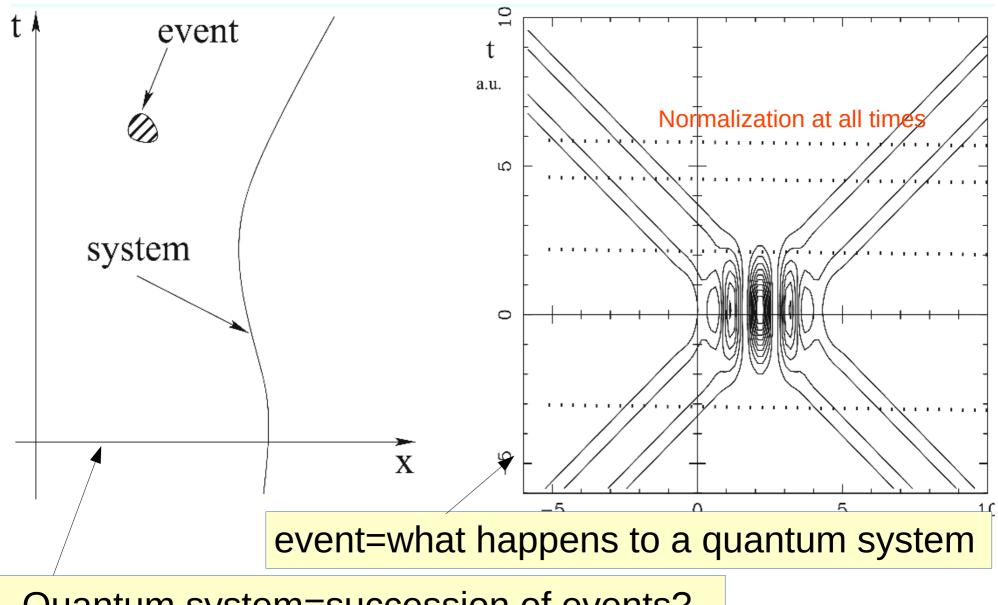
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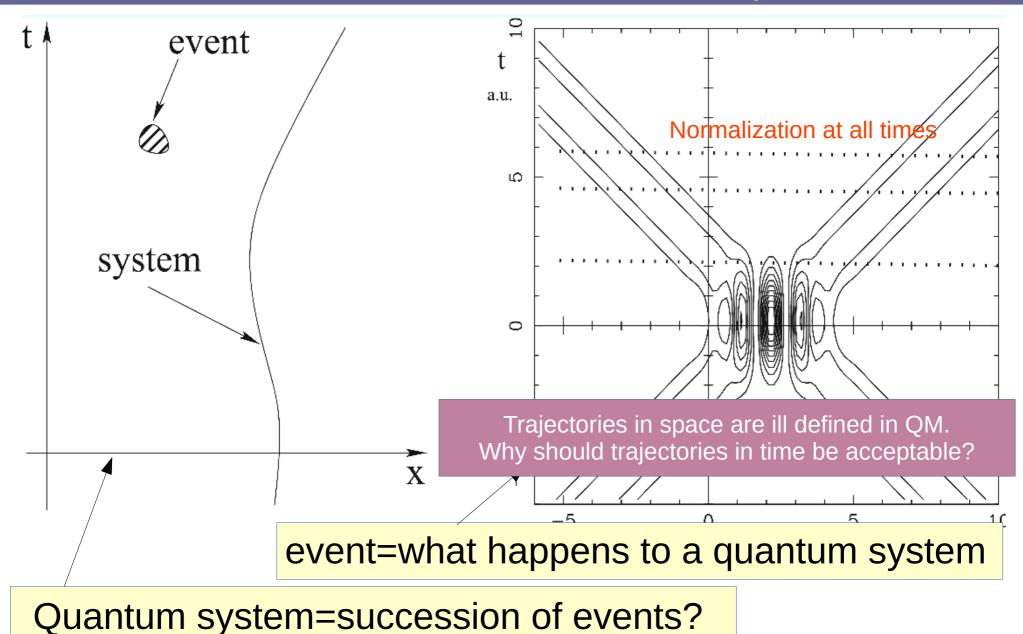
Explore the alternative!!

#### Current QM not able to deal with q events

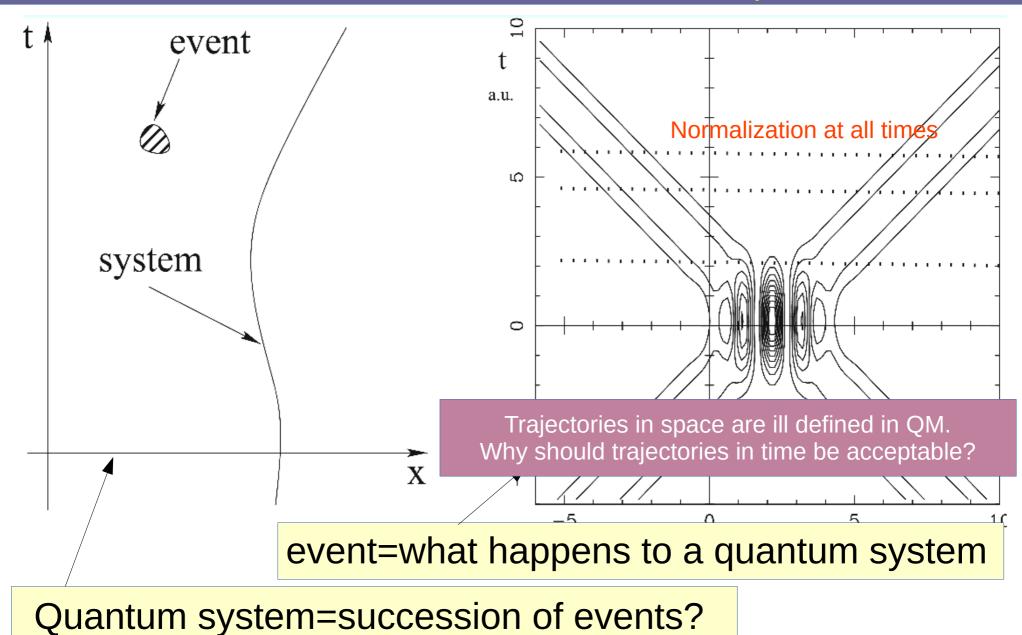


Quantum system=succession of events?

#### Current QM not able to deal with q events



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Need: Hilbert space for events (and its composition rule!)

# Start with SPECIAL relativity

Let's deal with GR in the future (much in the future!)

QM uses time conditioned quantities



**Textbook QM and QFT** 

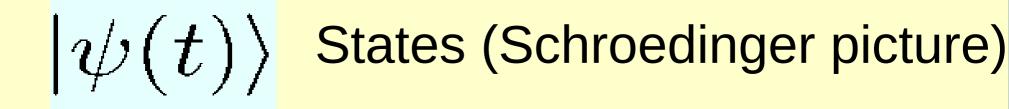
QM uses time conditioned quantities

# $|\psi(t) angle$ States (Schroedinger picture)X(t) Observables (Heis. picture)



**Textbook QM and QFT** 

QM uses time conditioned quantities



### X(t) Observables (Heis. picture)

### CANNOT be relativistically covariant (covariance="formulas look the same in all reference frames")

### Wait?!? What about QFT?



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- 1) use observables in the Heisenberg picture with covariant spacetime dependence, e.g.  $a^{\dagger} e^{-ix^{\mu}p_{\mu}}$



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2) Use a state that is invariant for Lorentz transforms, e.g the vacuum  $|0\rangle$ 





### Our approach: Geometric Event-Based QM







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### quantum events → fundamental



### Our approach: Geometric Event-Based QM





### quantum events → fundamental

quantum systems  $\rightarrow$  derived: a quantum state for a succession of events in q spacetime





(without energy nothing can happen, without momentum, nothing can be localized)



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### basic observables:

$$\overline{X} := (X^0, X^1, X^2, X^3) \overline{P} := (P^0, P^1, P^2, P^3)$$



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why?!?



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why?!? Poincare' algebra: $\begin{bmatrix} M^{\mu\nu}, P^{\rho} \end{bmatrix} = -i(\eta^{\mu\rho}P^{\nu} - \eta^{\nu\rho}P^{\mu}), \\ \begin{bmatrix} M^{\mu\nu}, M^{\rho\sigma} \end{bmatrix} = i(\eta^{\nu\rho}M^{\mu\sigma} - \eta^{\mu\rho}M^{\nu\sigma}) \\ -\eta^{\mu\sigma}M^{\rho\nu} + \eta^{\nu\sigma}M^{\rho\mu}) \end{bmatrix}$  Now we can do GEB of easy systems (scalar KG and Dirac). Can we extend to more complex fields?

Can we extend to GR?



$$|\Phi\rangle = \int d^4x \; \Phi(\overline{x}) \; |\overline{x}\rangle = \int d^4p \; \tilde{\Phi}(\overline{p}) \; |\overline{p}\rangle$$



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 $\left|\overline{x}\right\rangle$  is not (it describes a detection event of a particle that exists only at one spt location)

It cannot describe a quantum system (i.e. a succession of events), but it can describe a measurement apparatus.



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Physical states are the ones that satisfy some constraints (e.g. PW WdW).



$$|\Phi\rangle = \int d^4x \; \Phi(\overline{x}) \; |\overline{x}\rangle = \int d^4p \; \tilde{\Phi}(\overline{p}) \; |\overline{p}\rangle$$

Amplitudes (wavefunctions)

$$\Phi(\overline{x}) := \langle \overline{x} | \Phi \rangle , \ \tilde{\Phi}(\overline{p}) := \langle \overline{p} | \Phi \rangle = \int \frac{d^4x}{4\pi^2} e^{i\overline{x} \cdot \underline{p}} \ \Phi(\overline{x})$$



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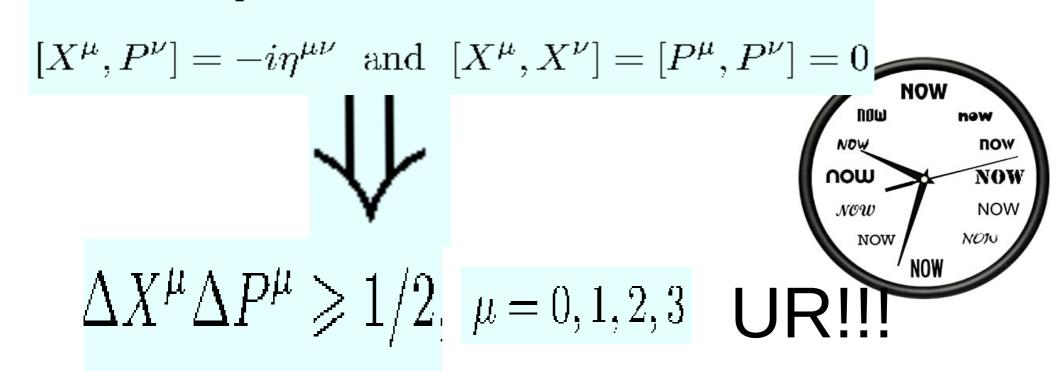
# **UNCONDITIONED** probability that the event is in spacetime position $\overline{x}$

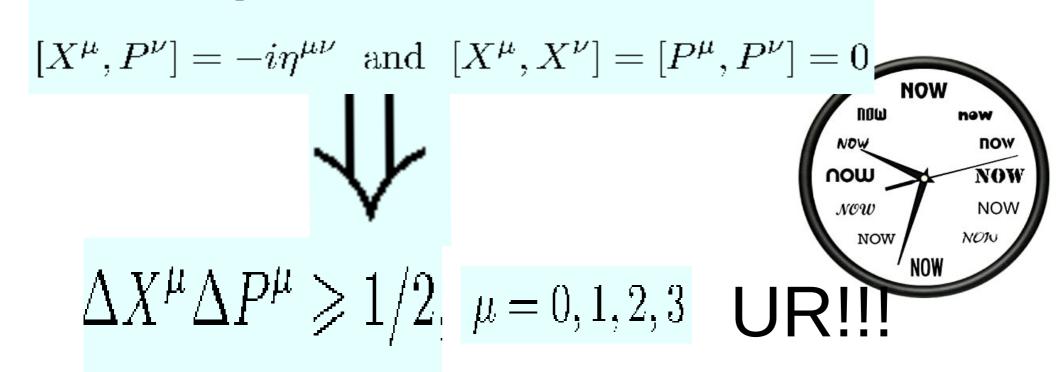
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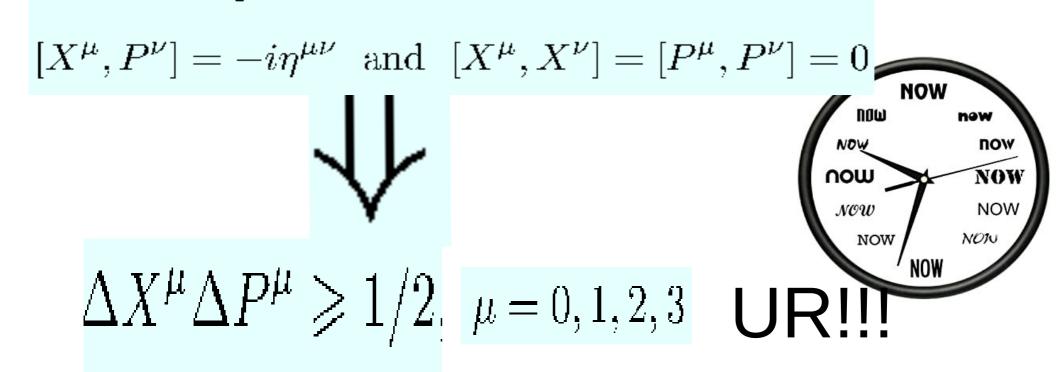
Born rule in QM is **CONDITIONED**   $p(\vec{x}|\psi,t) = |\langle \vec{x}_S | \psi_S(t) \rangle|^2 = [\langle \vec{x}_H(t) | \psi_H \rangle|^2$ Probability that the particle is at position  $\vec{x}$ *Given* that the time is t!!! QM

### Born rule: $P(\overline{x}|\Phi) = |\Phi(\overline{x})|^2 = |\langle \overline{x}|\Phi\rangle|^2$ **UNCONDITIONED** probability that the event is in spacetime position $\overline{x}$ QM probabilities are NOT covariant **GEB** probabilities ARE covariant Born rule in QM is **CONDITIONED** $p(\vec{x}|\psi,t) = |\langle \vec{x}_S |\psi_S(t)\rangle|^2 = |\langle \vec{x}_H(t)|\psi_H\rangle|^2$ Probability that the particle is at position $\vec{x}$ *QIVEN* that the time is t!!!





In GEB  $\Delta X^0 \Delta P^0 \ge \hbar/2$  is a Heisenberg-Robertson inequality, in QM it is completely meaningless (e.g. Peres, Aharonov-Bohm)



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Cannot localize an event in time unless it has an energy spread

### LORENTZ TRANSFORMS IN GEB

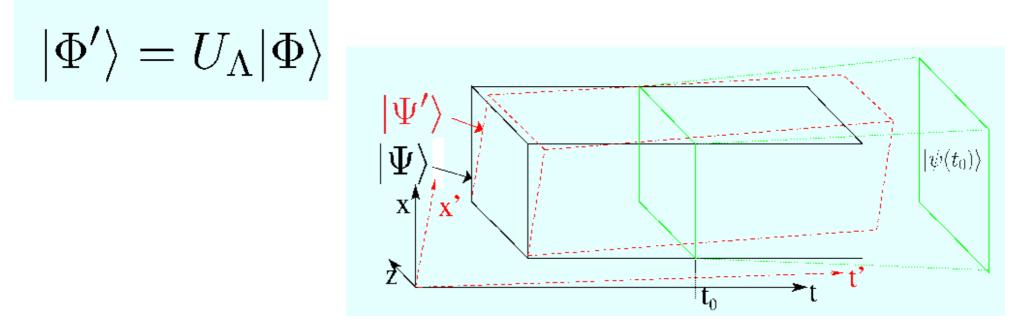
## LORENTZ TRANSFORMS IN GEB

Just a **unitary transformation on the GEB state** (Wigner's prescription on how to describe symmetries of a theory)

$$|\Phi'\rangle = U_{\Lambda}|\Phi\rangle$$

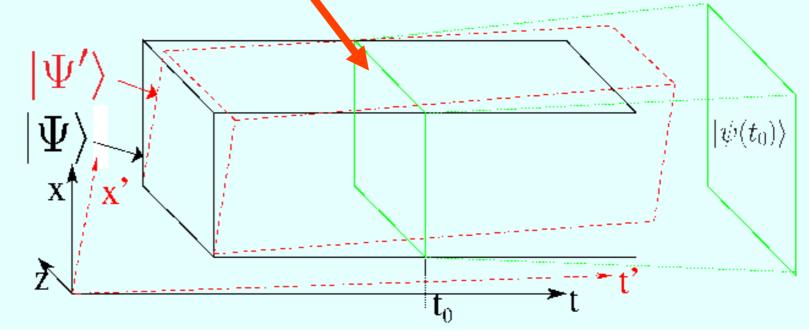
## LORENTZ TRANSFORMS IN GEB

Just a **unitary transformation on the GEB state** (Wigner's prescription on how to describe symmetries of a theory)



 $|\Phi\rangle = \int d^4x \; \Phi(\overline{x}) \; |\overline{x}\rangle = \int d^4p \; \tilde{\Phi}(\overline{p}) \; |\overline{p}\rangle$ 

#### 

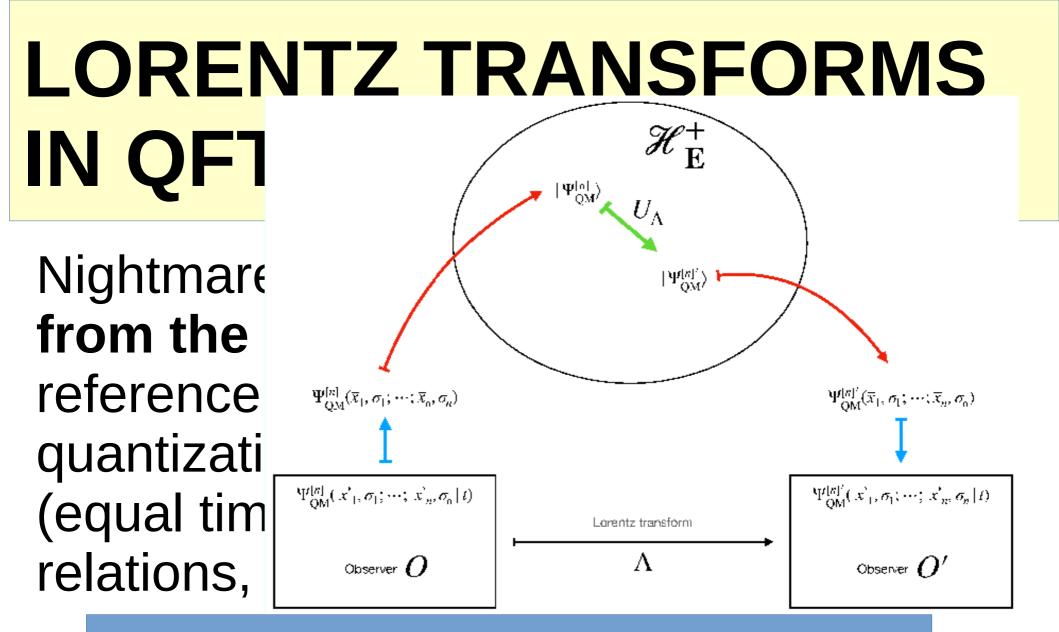


Nightmare!

Nightmare! Need to **quantize from the start** in the new reference frame: rerun the quantization procedure (equal time commutation relations, etc.).

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... or you can take a shortcut through GEB



Easier than requantizing everything: a good first motivation for GEB

GEB

(if fixed number of events n)

$$|\Phi^{[n]}\rangle = \sum_{\sigma_1,\cdots,\sigma_n} \int d^4 x_1 \cdots d^4 x_n \; \Phi^{[n]}(\overline{x}_1,\sigma_1;\cdots;\overline{x}_n,\sigma_n) \Big| \overline{x}_1,\sigma_1;\cdots;\overline{x}_n,\sigma_n \Big\rangle$$



(if fixed number of events n)

Fock space

$$|\Phi^{[n]}\rangle = \sum_{\sigma_1,\cdots,\sigma_n} \int d^4 x_1 \cdots d^4 x_n \; \Phi^{[n]}(\overline{x}_1,\sigma_1;\cdots;\overline{x}_n,\sigma_n) \Big| \overline{x}_1,\sigma_1;\cdots;\overline{x}_n,\sigma_n \big\rangle$$

#### (otherwise)

$$=\frac{1}{\sqrt{n!}}\sum_{\sigma_1,\cdots,\sigma_n}\int d^4x_1\cdots d^4x_n \,\Phi^{[n]}(\overline{x}_1,\sigma_1;\cdots;\overline{x}_n,\sigma_n) \,a\frac{\dagger}{\overline{x}_1,\sigma_1}\cdots a\frac{\dagger}{\overline{x}_n,\sigma_n}|0\rangle_4$$

## Creation operators: create an event at position $x_1$

(if fixed number of events n)

$$\begin{split} |\Phi^{[n]}\rangle &= \sum_{\sigma_{1},\cdots,\sigma_{n}} \int d^{4}x_{1}\cdots d^{4}x_{n} \Phi^{[n]}(\overline{x}_{1},\sigma_{1};\cdots;\overline{x}_{n},\sigma_{n}) |\overline{x}_{1},\sigma_{1};\cdots;\overline{x}_{n},\sigma_{n}\rangle \\ \hline \textbf{Fock space} \quad \textbf{(otherwise)} \\ &= \frac{1}{\sqrt{n!}} \sum_{\sigma_{1},\cdots,\sigma_{n}} \int d^{4}x_{1}\cdots d^{4}x_{n} \Phi^{[n]}(\overline{x}_{1},\sigma_{1};\cdots;\overline{x}_{n},\sigma_{n}) |u_{\overline{x}_{1},\sigma_{1}}^{\dagger}\cdots u_{\overline{x}_{n},\sigma_{n}}^{\dagger}|0\rangle_{4} \\ P^{[n]}(\overline{x}_{1},\sigma_{1};\cdots;\overline{x}_{n},\sigma_{n}) &= |\Phi^{[n]}(\overline{x}_{1},\sigma_{1};\cdots;\overline{x}_{n},\sigma_{n})|^{2} \end{split}$$

Joint probability for the n events to happen in spt positions  $x_1...x_n$ 

(if fixed number of events n)

 $|\Phi^{[n]}
angle = \sum \int d^4x_1 \cdots d^4x_n \Phi^{[n]}(\overline{x}_1, \sigma_1; \cdots; \overline{x}_n, \sigma_n) |\overline{x}_1, \sigma_1; \cdots; \overline{x}_n, \sigma_n
angle$ 

### EACH EVENT WITH ITS OWN TIME!!!! (cfr Dirac's multiparticle-multitime)

$$P^{[n]}(\overline{x}_1, \sigma_1; \cdots; \overline{x}_n, \sigma_n) = |\Phi^{[n]}(\overline{x}_1, \sigma_1; \cdots; \overline{x}_n, \sigma_n)|^2$$

Joint probability for the n events to happen in spt positions  $x_1...x_n$ 

## $=\frac{1}{\sqrt{n!}}\sum_{\sigma_1,\cdots,\sigma_n}\int d^4x_1\cdots d^4x_n \,\Phi^{[n]}(\overline{x}_1,\sigma_1;\cdots;\overline{x}_n,\sigma_n) \,\left[a\frac{\dagger}{\overline{x}_1},\sigma_1\cdots a\frac{\dagger}{\overline{x}_n},\sigma_n\right]|0\rangle_4$



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### Commutators:

Bose: 
$$[a_{\overline{x},\sigma}, a_{\overline{x}',\sigma'}^{\dagger}] = \delta_{\sigma,\sigma'} \delta^{(4)}(\overline{x} - \overline{x}'), \ [a_{\overline{x},\sigma}, a_{\overline{x}',\sigma'}] = 0,$$
  
Fermi:  $\{a_{\overline{x},\sigma}, a_{\overline{x}',\sigma'}^{\dagger}\} = \delta_{\sigma,\sigma'} \delta^{(4)}(\overline{x} - \overline{x}'), \ \{a_{\overline{x},\sigma}, a_{\overline{x}',\sigma'}\} = 0$ 



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Bosonic events — Bosons

Fermionic events — Fermions

9-1 6 V

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4D vacuum



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#### 4D vacuum

# **Different** state from the 3D vacuum of QFT $(|0\rangle_3 = n0)$ particles at time t=0 in the Heis pic)



$$=\frac{1}{\sqrt{n!}}\sum_{\sigma_1,\cdots,\sigma_n}\int d^4x_1\cdots d^4x_n \,\Phi^{[n]}(\overline{x}_1,\sigma_1;\cdots;\overline{x}_n,\sigma_n) \,a\frac{\dagger}{\overline{x}_1,\sigma_1}\cdots a\frac{\dagger}{\overline{x}_n,\sigma_n} |0\rangle_4$$

### 

vacuum

$$|0\rangle_3 = \text{foliate}(a_{\overline{p}=0}^{\dagger}|0\rangle_4)$$

Event state of **zero 4-momentum**: ground state of the field

#### Relativistic QM FROM GEB



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# Use **constraints** (as in the quantum time P&W, WdW, etc.)!



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(GEB state describing the whole dynamics of the **particle as a state of a sequence of events**)



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 $K|\Psi_{\rm QM}\rangle = 0$ 

(GEB state describing the whole dynamics of the **particle as a state of a sequence of events**)

 Write it as an eigenstate of a constraint op.



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Same procedure works for more complex QFT systems e.g.



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Similarly for the Dirac eq. constraint.

No claim that QFT is incorrect.

GEB is an **alternative** in the sense "a different way to obtain the same results"

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WHY?

No claim that QFT is incorrect.

# GEB is an **alternative** in the sense "a different way to obtain the same results"

WHY?

To get a better ontology? To go further than QFT can go?

## Conclusions

• Philosophical considerations about time



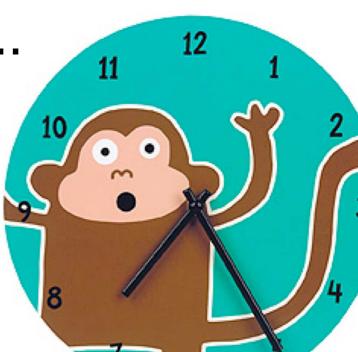
- Philosophical considerations about time
- •Time as a quantum degree of freedom



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11

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11

- Pauli objections and others..
- Quantum time measurements
- GEB a relativistically covariant quantiz

Take home message

## A quantization of time based on conditional probability amplitudes



quantum time: PRD **92**, 045033 Geometric Event-Based QM: NJP **25**, 023027 time observable: PRL **124**, 110402

Lorenzo Maccone maccone@unipv.it