Minimum Detection Efficiencies for Genuine Nonlocality Tests

ICQIST-2025

Subhendu Bikash Ghosh

SNBNCBS, Kolkata

• Bell's theorem says Quantum theory is nonlocal.

• Is nature nonlocal?

What Schrödinger found problematic—indeed, objectionable—about entanglement was this possibility of remote steering (Ref. 43, p. 556):

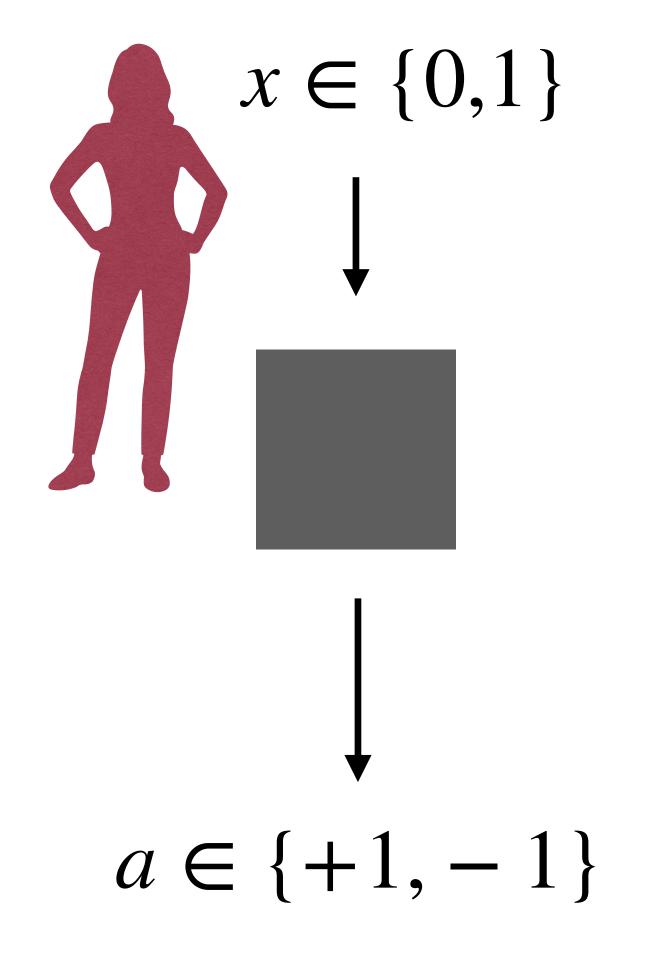
It is rather discomforting that the theory should allow a system to be steered or piloted into one or the other type of state at the experimenter's mercy in spite of his having no access to it.

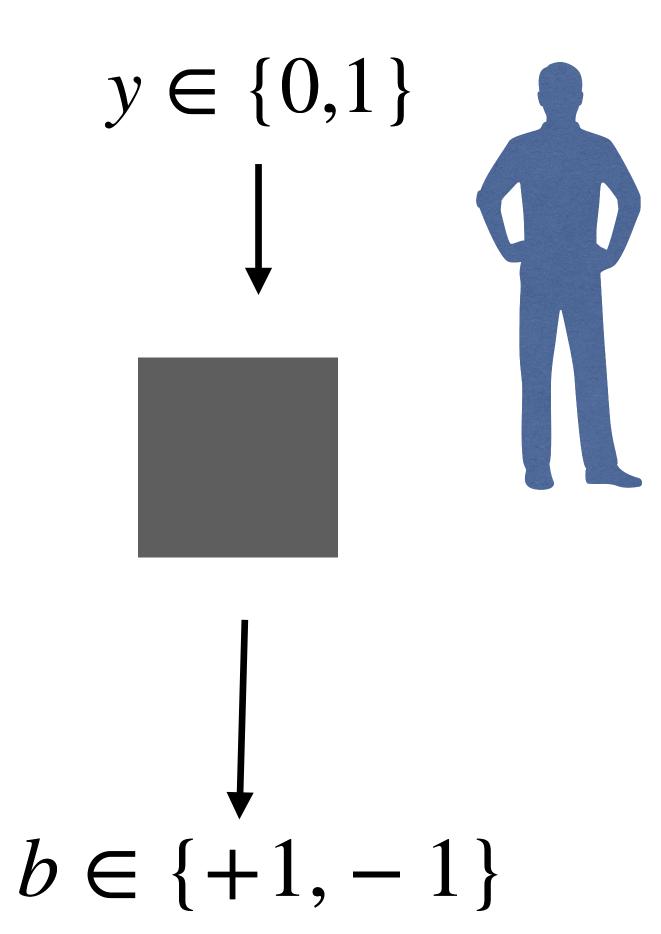
R. Clifton et al. Foundations of Physics 33, 1561-1591 (2003)

He conjectured that an entangled state of a composite system would almost instantaneously decay to a mixture as the component systems separated. (A similar possibility was raised and rejected by Furry⁽²⁶⁾.) There would still be correlations between the states of the component systems, but remote steering would no longer be possible (Ref. 44, p. 451):

It seems worth noticing that the [EPR] paradox could be avoided by a very simple assumption, namely if the situation after separating were described by the expansion (12), but with the additional statement that the knowledge of the *phase relations* between the complex constants a_k has been entirely lost in consequence of the process of separation. This would mean that not only the parts, but the whole system, would be in the situation of a mixture, not of a pure state. It would not preclude the possibility of determining the state of the first system by *suitable* measurements in the second one or *vice versa*. But it would utterly eliminate the experimenters influence on the state of that system which he does not touch.

R. Clifton et al. Foundations of Physics 33, 1561-1591 (2003)



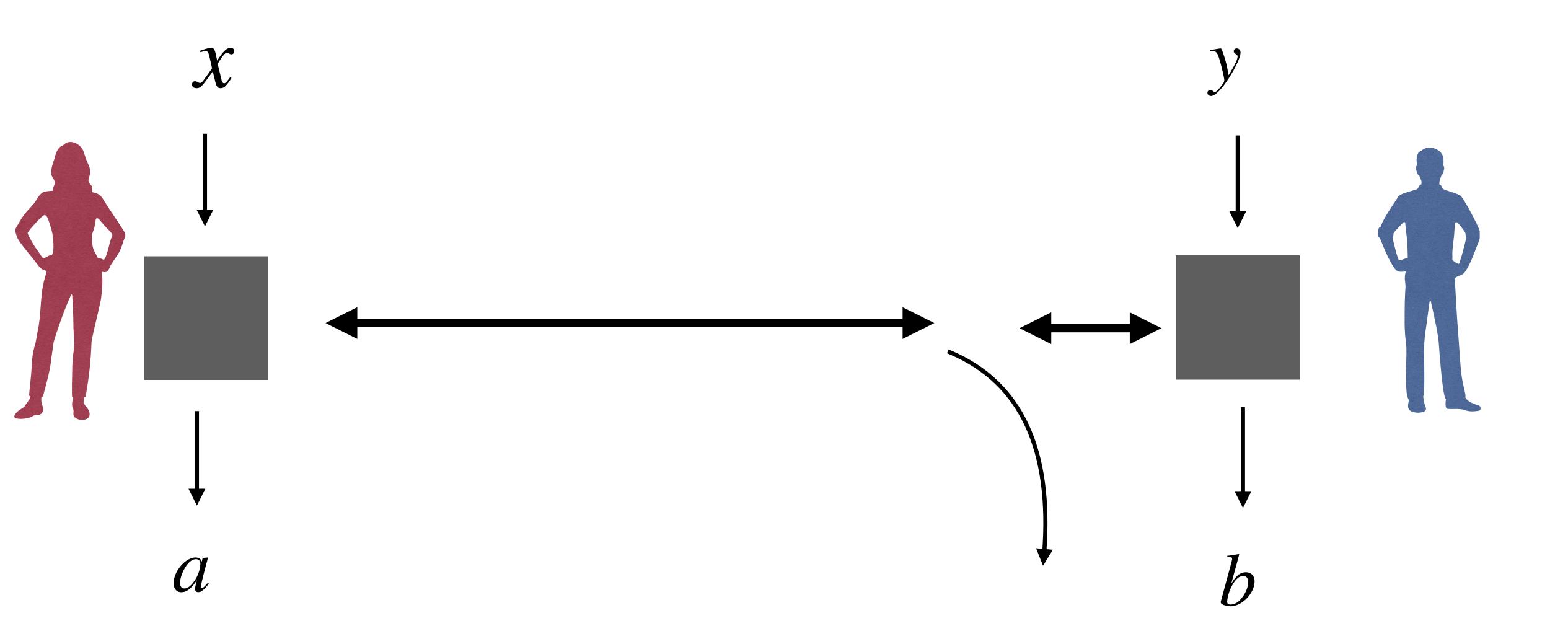


- Locality loophole
- Detection loophole
- Free will loophole

Locality loophole

• Detection loophole

• Free will loophole



 Danger of post selection (P.M. Pearl Phys. Rev. D 2, 1418(1970)

- Danger of post selection (P.M. Pearl Phys. Rev. D 2, 1418(1970)
- No click events as valid outcome.

- Danger of post selection (P.M. Pearl Phys. Rev. D 2, 1418(1970)
- No click events as valid outcome.
- Non maximally entangled states are more efficient for a non efficient detector. (Jan-Åke Larsson 2014 J. Phys. A: Math. Theor. 47 424003)

- Danger of post selection (P.M. Pearl Phys. Rev. D 2, 1418(1970)
- No click events as valid outcome.
- Non maximally entangled states are more efficient for a non efficient detector. (Jan-Åke Larsson 2014 J. Phys. A: Math. Theor. 47 424003)
- We ask the question for genuine nonlocality.

 Danger of post selection (P.M. Pearl Phys. Rev. D 2, 1418(1970)

No click PHYSICAL REVIEW A covering atomic, molecular, and optical physics and quantum information Non max for a Letters Recent Accepted Collections Authors Referees Search About **Editorial Team** non effic hys. A: Math. Th Minimum detection efficiencies for loopholefree genuine nonlocality tests Subhendu B. Ghosh, Snehasish Roy Chowdhury, Ranendu Adhikary, Arup Roy, and Tamal Guha We ask Phys. Rev. A 109, 052202 – Published 2 May 2024 ₩ More

T_2 nonlocality

First we consider a operationally motivated class of genuine nonlocality, T_2 nonlocality. (Phys. Rev. A 88, 014102 (2013))

T_2 nonlocality

First we consider a operationally motivated class of genuine nonlocality, T_2 nonlocality. (Phys. Rev. A 88, 014102 (2013))

A tripartite correlation $P(abc \mid xyz)$ will be T_2 nonlocal if it violates the following inequality.

$$\begin{split} \mathcal{B}_{T_2} &= -2\{P(00\,|A_1B_1) + P(00\,|B_1C_1) + P(00\,|C_1A_1)\} \\ &- P(000\,|A_0B_0C_1) - P(000\,|A_0B_1C_0) - P(000\,|A_1B_0C_0) \\ &+ 2\{P(000\,|A_0B_1C_1) + P(000\,|A_1B_0C_1) + P(000\,|A_1B_1C_0) + P(000\,|A_1B_1C_1\} \leq 0 \end{split}$$

- We assign no click events as '1'.
- ullet Three parties with detectors of efficiencies η_A, η_B and η_C .

• We assign no click events as '1'.

Lemma 1: The spatially separated parties would be able to certify genuine three-way nonlocality in terms of the inequality \mathcal{B}_{T_2} , when the following condition holds

$$(4\eta_A \eta_B \eta_C - \eta_A \eta_B - \eta_A \eta_C - \eta_B \eta_C) > 0$$

This necessary condition holds for any no-signaling theory.

Lemma 2: There exists a quantum settings with a single parameter θ , which violates the inequality

 \mathcal{B}_{T_2} , whenever $(4\eta_A\eta_B\eta_C-\eta_A\eta_B-\eta_A\eta_C-\eta_B\eta_C)>0$.

Lemma 2: There exists a quantum settings with a single parameter θ , which violates the inequality

$$\mathcal{B}_{T_2}$$
, whenever $(4\eta_A\eta_B\eta_C-\eta_A\eta_B-\eta_A\eta_C-\eta_B\eta_C)>0$.

$$|\psi_{ABC}(\theta)\rangle = k(\theta)[(|011\rangle + |101\rangle + |110\rangle) + \frac{(1 - 3\cos\theta)}{\sin\theta}|111\rangle]$$

 $X_0 \equiv \{|0\rangle, |1\rangle\}$

 $X_1 = \{cos\theta | 0 > + sin\theta | 1 > , sin\theta | 0 > - cos\theta | 1 > \}$

Theorem 1: The minimum detection efficiencies for \mathcal{B}_{T_2} for each party must satisfy the following relation

$$(4\eta_A \eta_B \eta_C - \eta_A \eta_B - \eta_A \eta_C - \eta_B \eta_C) > 0$$

Theorem 1: The minimum detection efficiencies for \mathcal{B}_{T_2} for each party must satisfy the following relation

$$(4\eta_A \eta_B \eta_C - \eta_A \eta_B - \eta_A \eta_C - \eta_B \eta_C) > 0$$

In the case of $\eta_A=\eta_B=\eta_C=\eta$, the

$$\mathscr{B}_{T_2}$$
 will be violated iff $\eta > \frac{3}{4}$

- Previous result: MDE for Svetlitchny nonlocality is more than 97%. (V. Gebhart, PRA 106, 062202 (2022)).
- Our result shows 88.1% is sufficient.

Collaborators

- Snehasish, ISI
- Tamalda, Slovak Academy of Sciences
- Arupda, Hooghly Mohsin College
- Ranendu, ISI

Thank you

Table 1.2 A strategy exploiting the detection loophole for the CHSH test. In each round, Alice and Bob choose one of the eight quadruples of pre-established values that would give s = +2 (c.f., Table 1.1). Alice answers always, while Bob declines to answer to the input indicated by brackets, in such a way that the problematic output (boldface) is never produced.

$a_0, a_1; b_0, b_1$	a_0b_0	a_0b_1	a_1b_0	a_1b_1
+1, +1; +1, [+1]	+1	N	+1	N
-1, -1; -1, [-1]	+1	N	+1	N
+1, +1; +1, [-1]	+1	N	+1	N
-1, -1; -1, [+1]	+1	N	+1	N
+1, -1; [+1], +1	N	+1	\mathbf{N}	-1
-1, +1; [-1], -1	N	+1	\mathbf{N}	-1
+1, -1; [-1], +1	\mathbf{N}	+1	N	-1
-1, +1; [+1], -1	\mathbf{N}	+1	N	-1

Bell nonlocality, V. Scarani